

CMB Probes of LSS: Lensing & SZ

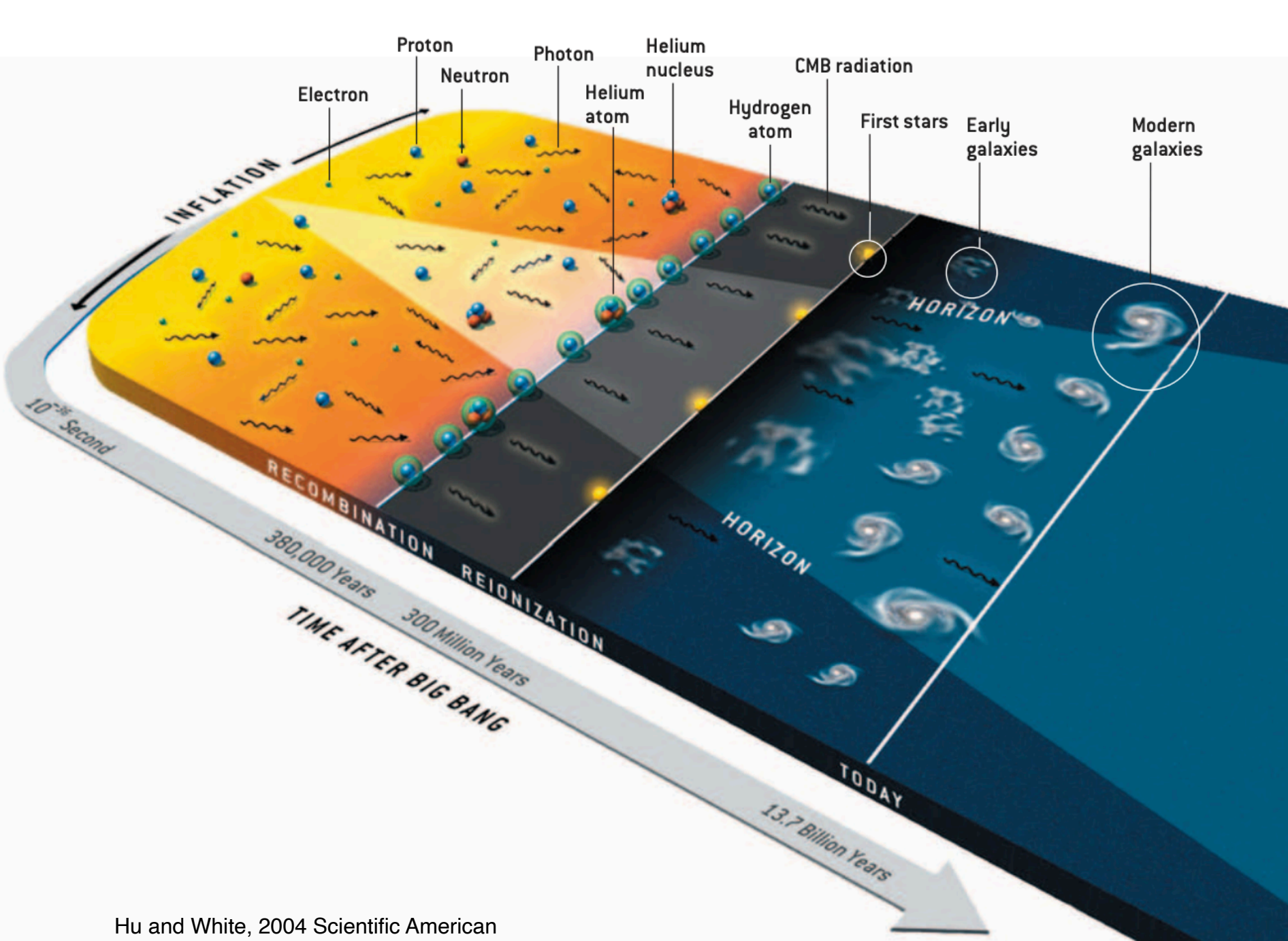
Gil Holder



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Outline

- the “surface of last scattering” is actually not the final word for lots of photons
 - ★ Thomson scattering
 - ★ lensing
 - ★ extragalactic foregrounds



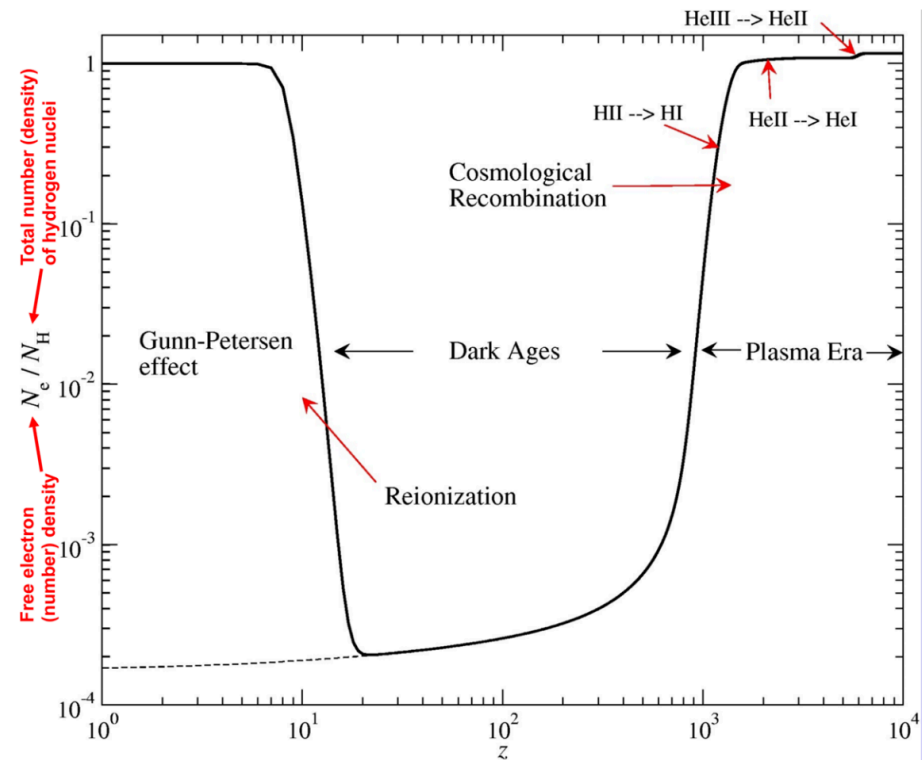
Hu and White, 2004 Scientific American

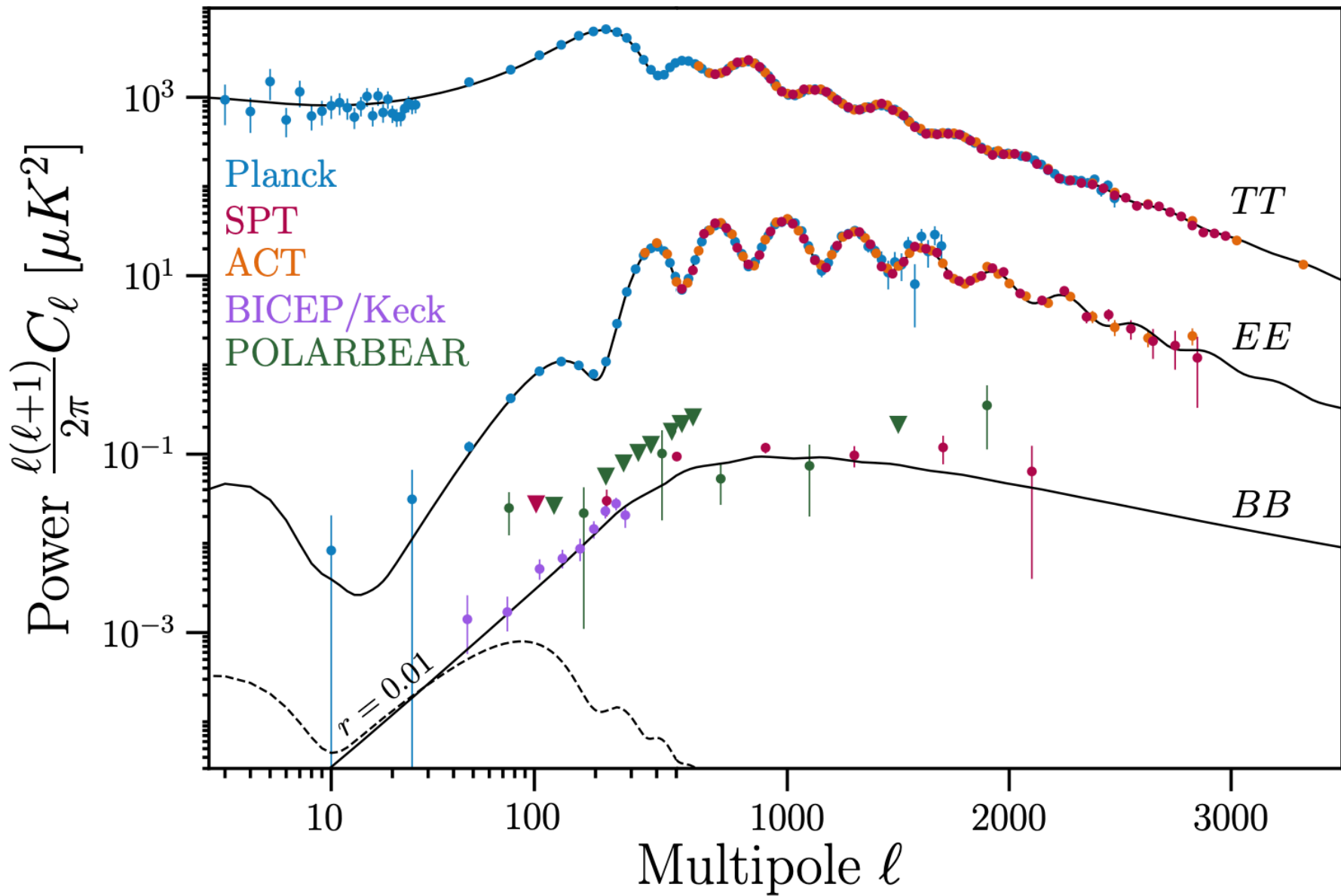
Ionization non-equilibrium

Hubble expansion causes recombinations to “freeze out” as e^- and p^+ can’t find each other in the dilute universe

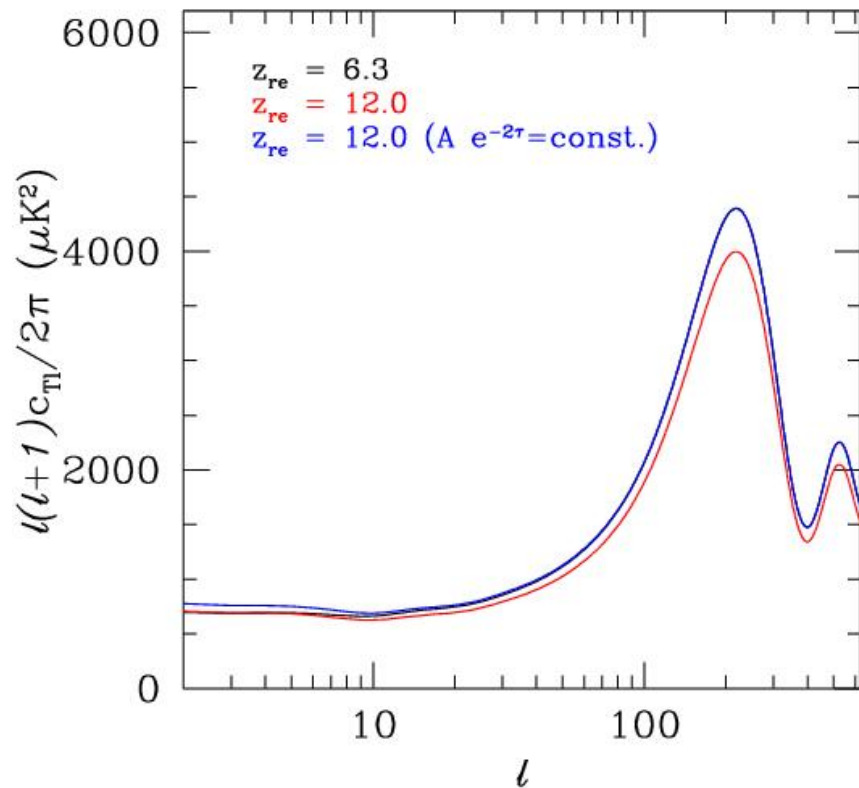
small residual ionization keeps gas and CMB thermally coupled for a surprisingly long time

reionization leads to unbinding of electrons from H atoms due to UV background ionizing field

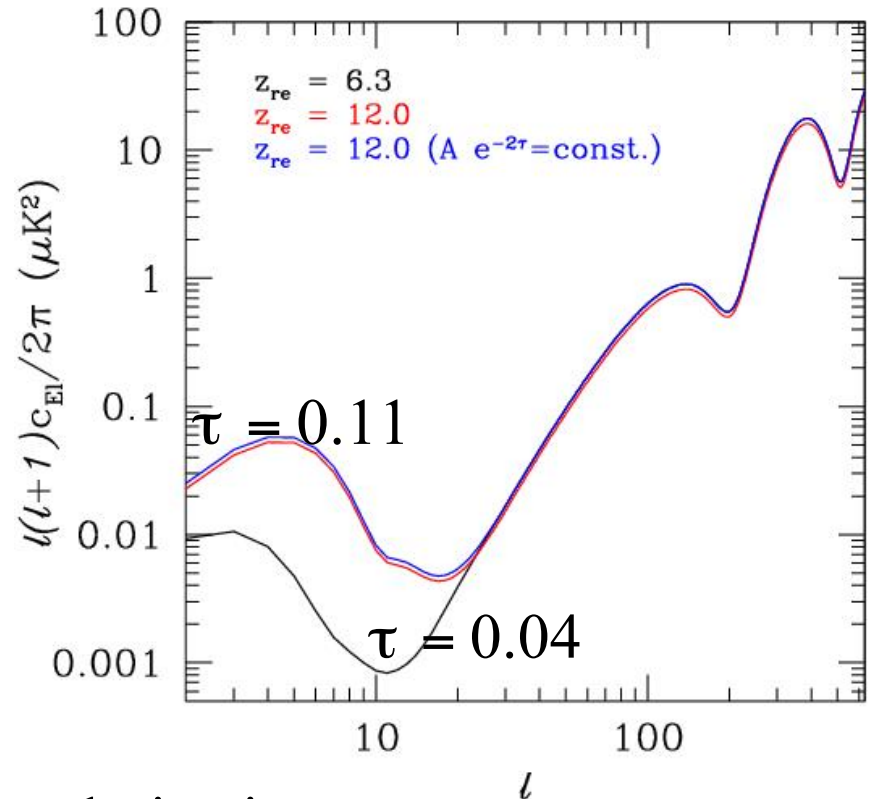




WMAP: ± 0.015 ; Planck: ± 0.005 ; ??? : ± 0.002



temperature

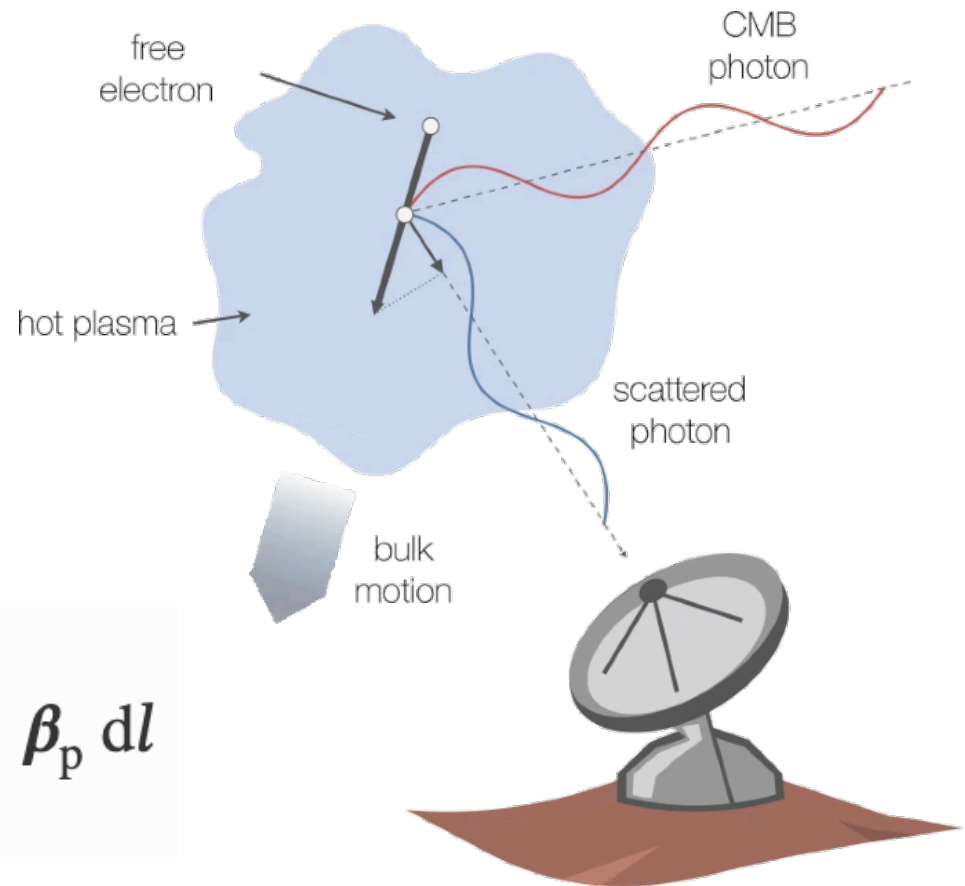


polarization

Scattering on moving electrons: kSZ

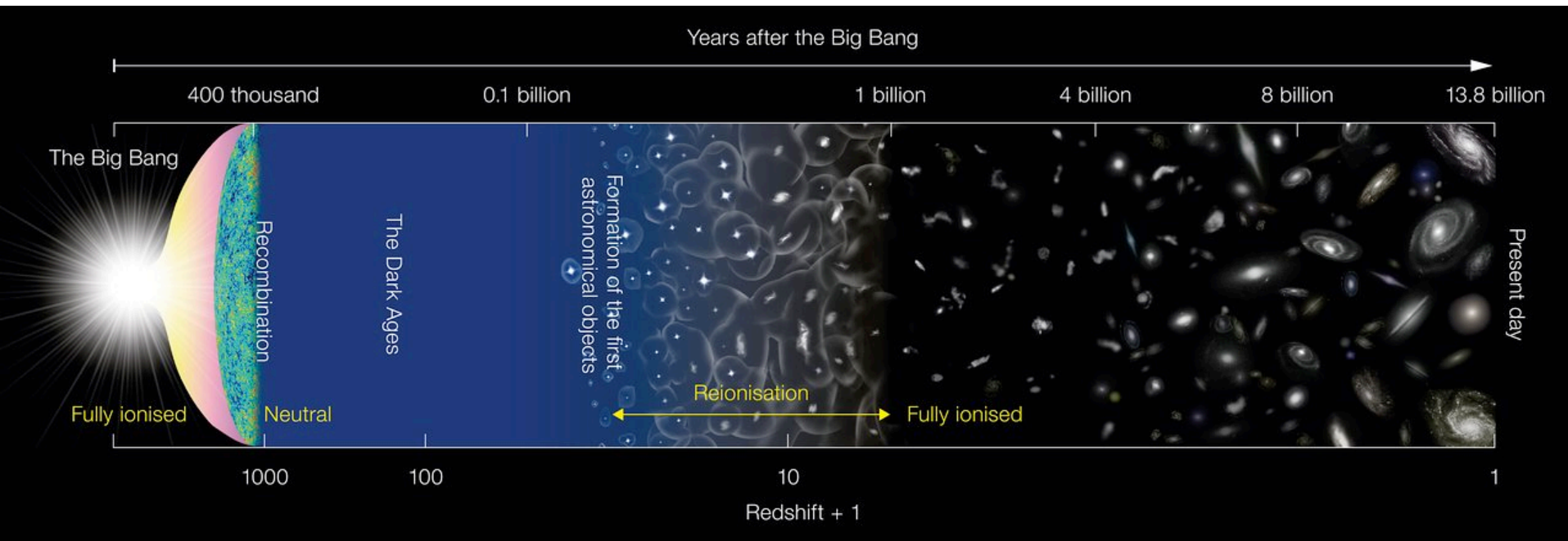
- kinetic Sunyaev-Zeldovich effect: Thomson scattering by bulk flow of electrons

$$\frac{\Delta T_{\text{CMB}}}{T_{\text{CMB}}} \approx - \int \sigma_{\text{T}} n_e n \cdot \beta_p dl$$

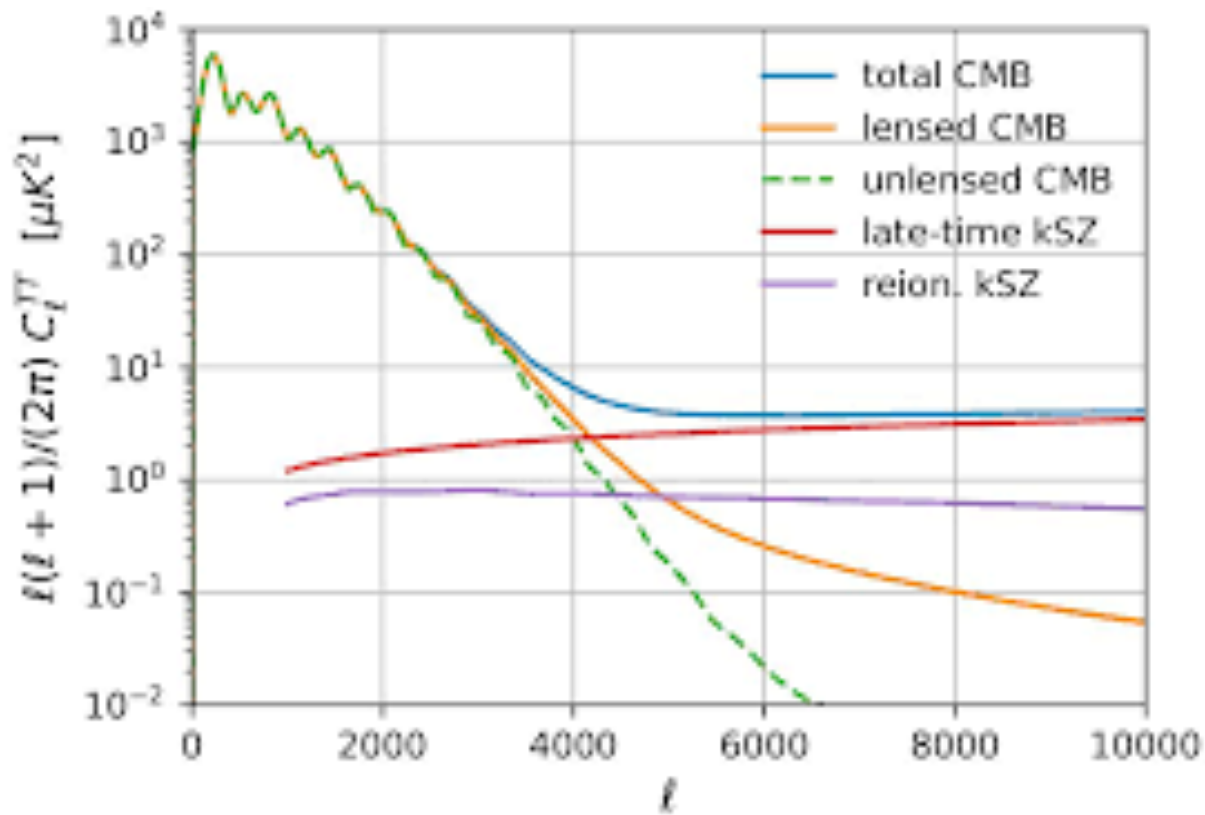


Scattering on moving electrons: kSZ

clumps of moving electrons at reionization, and at late times



Scattering on moving electrons: kSZ



Madhavacheril

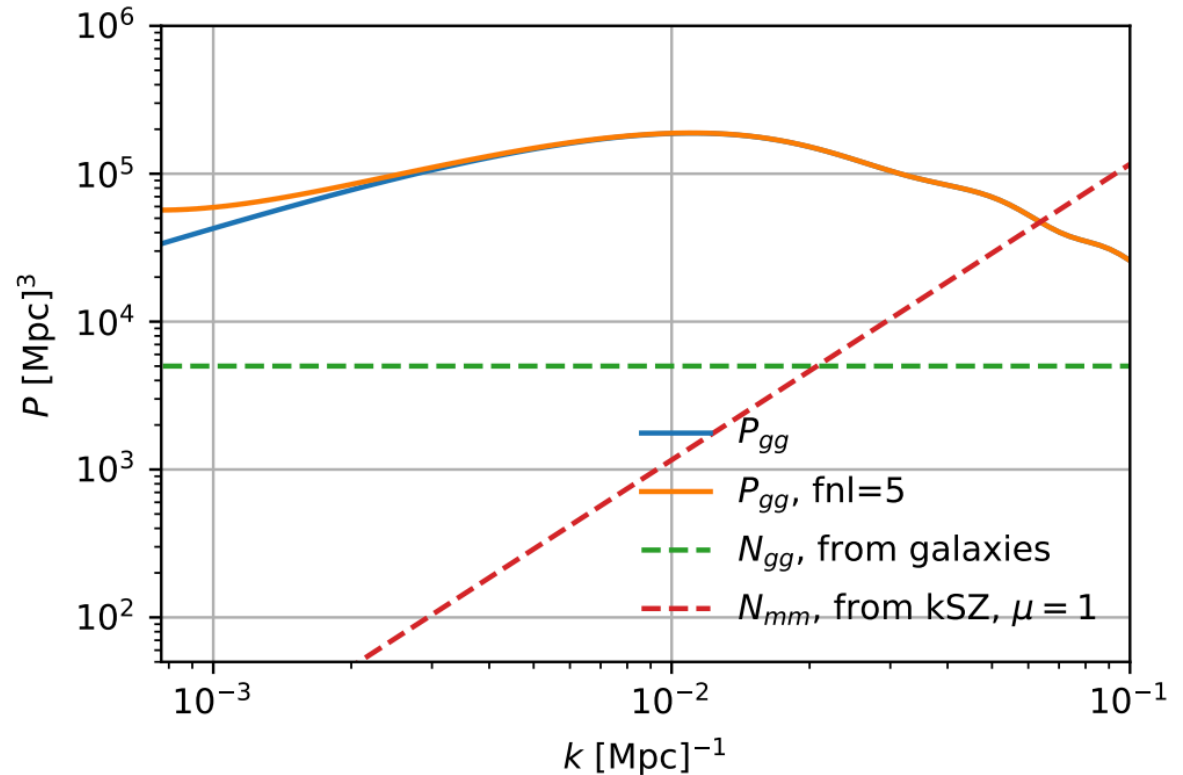
Scattering on moving electrons: kSZ

current status:

detected in cross-correlation with galaxies/clusters

forecast:

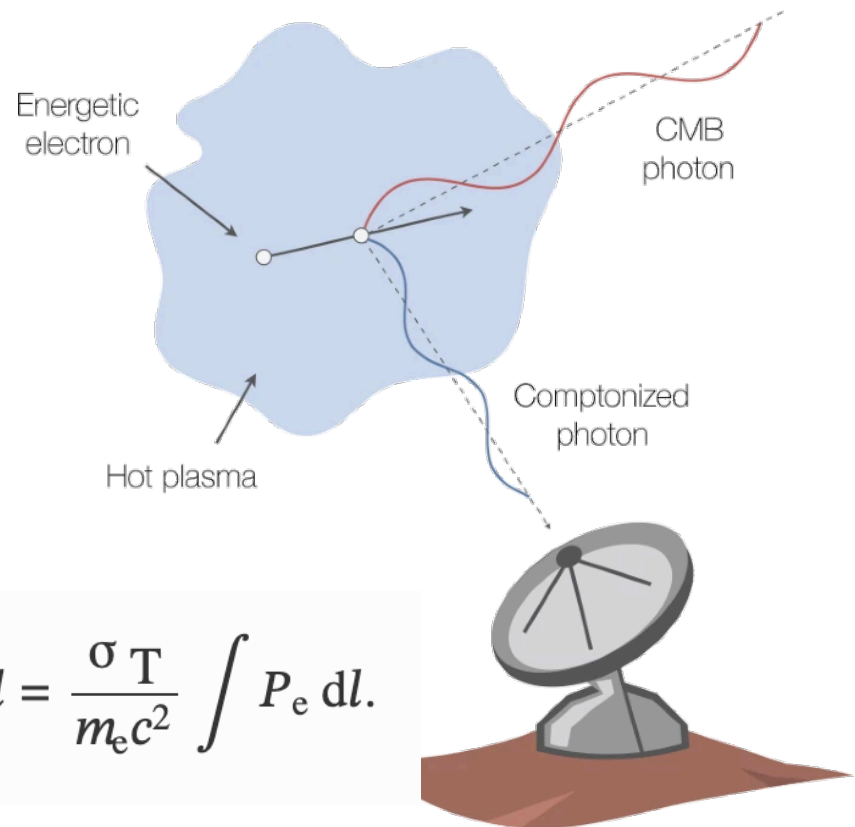
soon to be detected in auto-spectrum, higher order correlations could be very powerful for largest scales



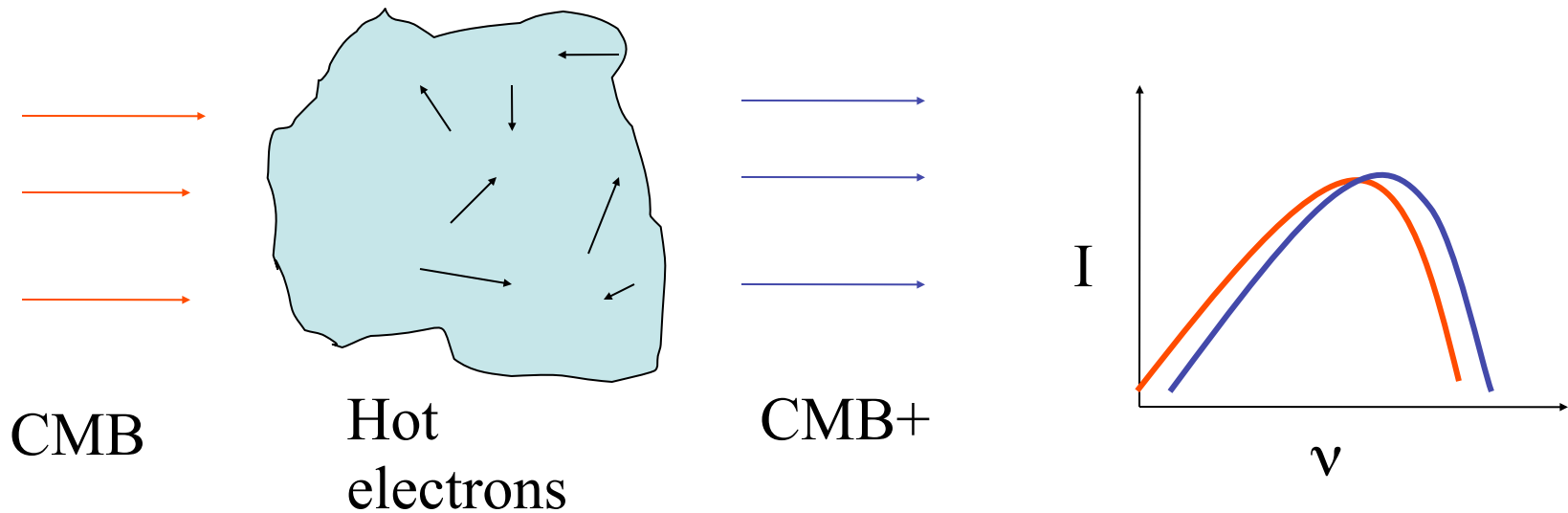
Scattering on moving electrons: tSZ

- thermal Sunyaev-Zeldovich effect: Thomson scattering by thermal motions of electrons

$$y \equiv \int \frac{k_B T_e}{m_e c^2} d\tau_e = \int \frac{k_B T_e}{m_e c^2} n_e \sigma_T dl = \frac{\sigma_T}{m_e c^2} \int P_e dl.$$



Thermal Sunyaev-Zel'dovich Effect



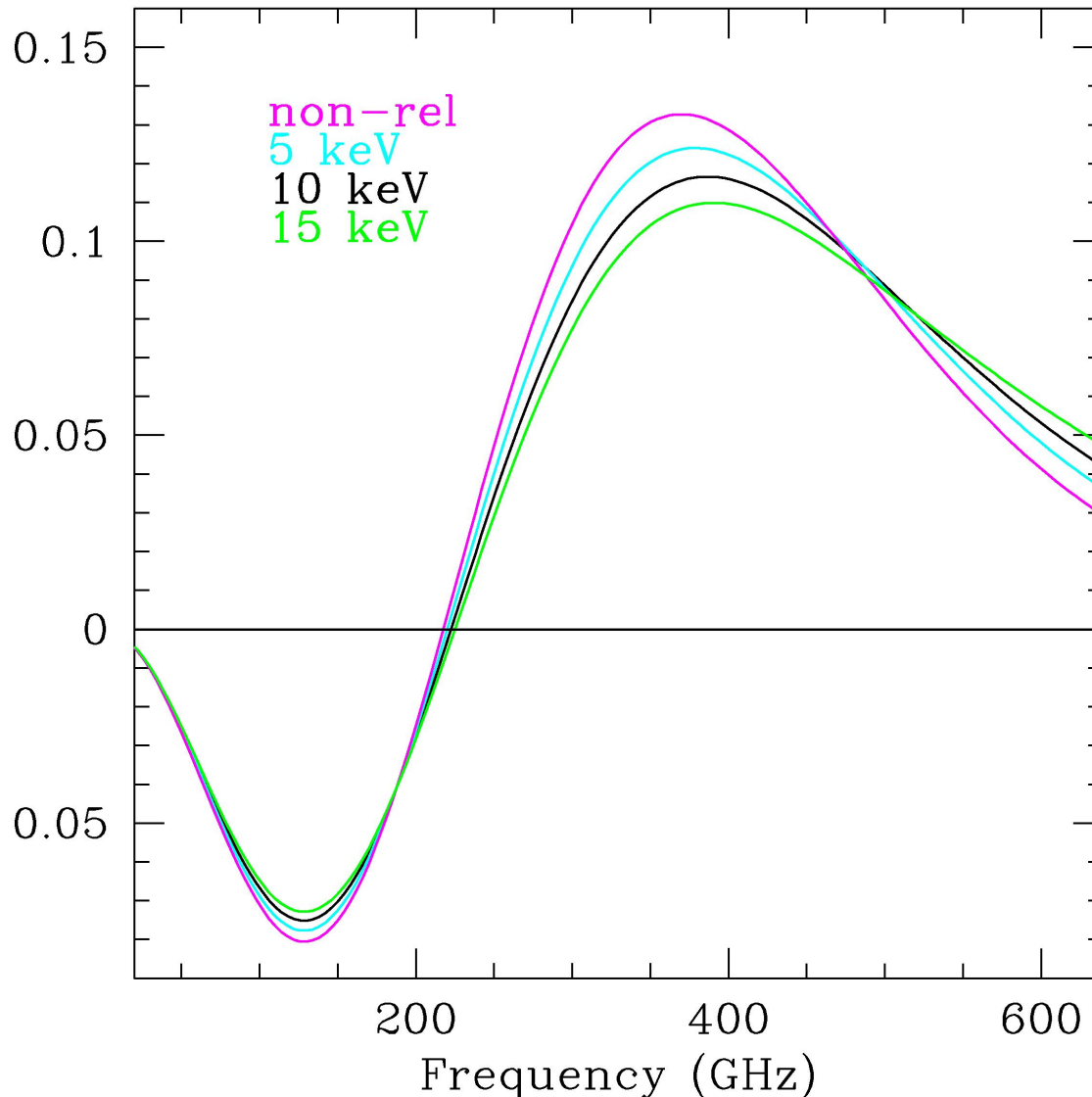
Optical depth: $\tau \sim 0.01$

Fractional energy gain per scatter: $\frac{kT}{m_e c^2} \sim 0.01$

Typical cluster signal: $\sim 500 \mu\text{K}$

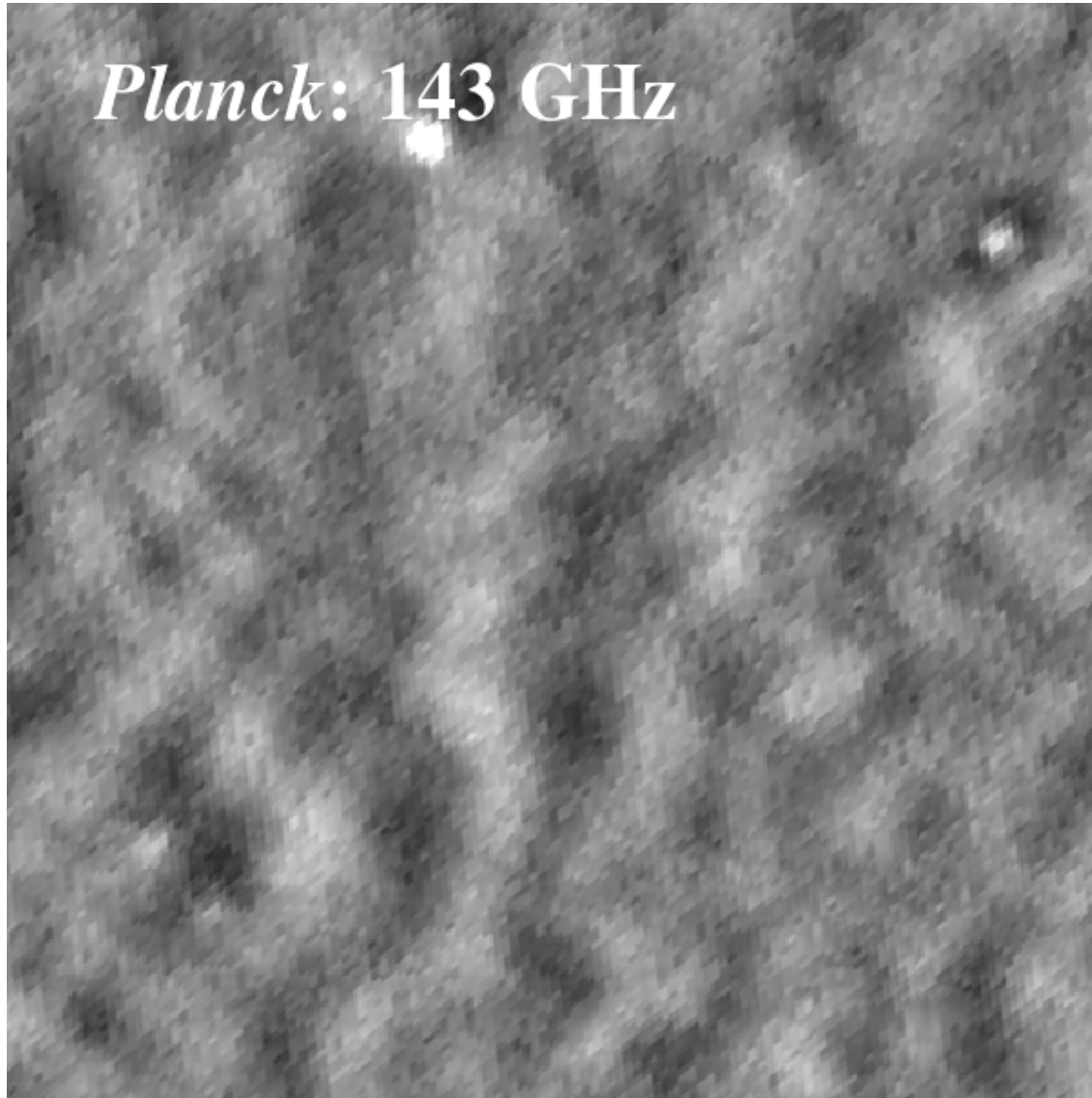
Thermal SZ Effect (and relativistic corrections)

uK imaging
would allow
1 keV
accuracy in
SZ
temperature



Planck: 143 GHz

3 degrees

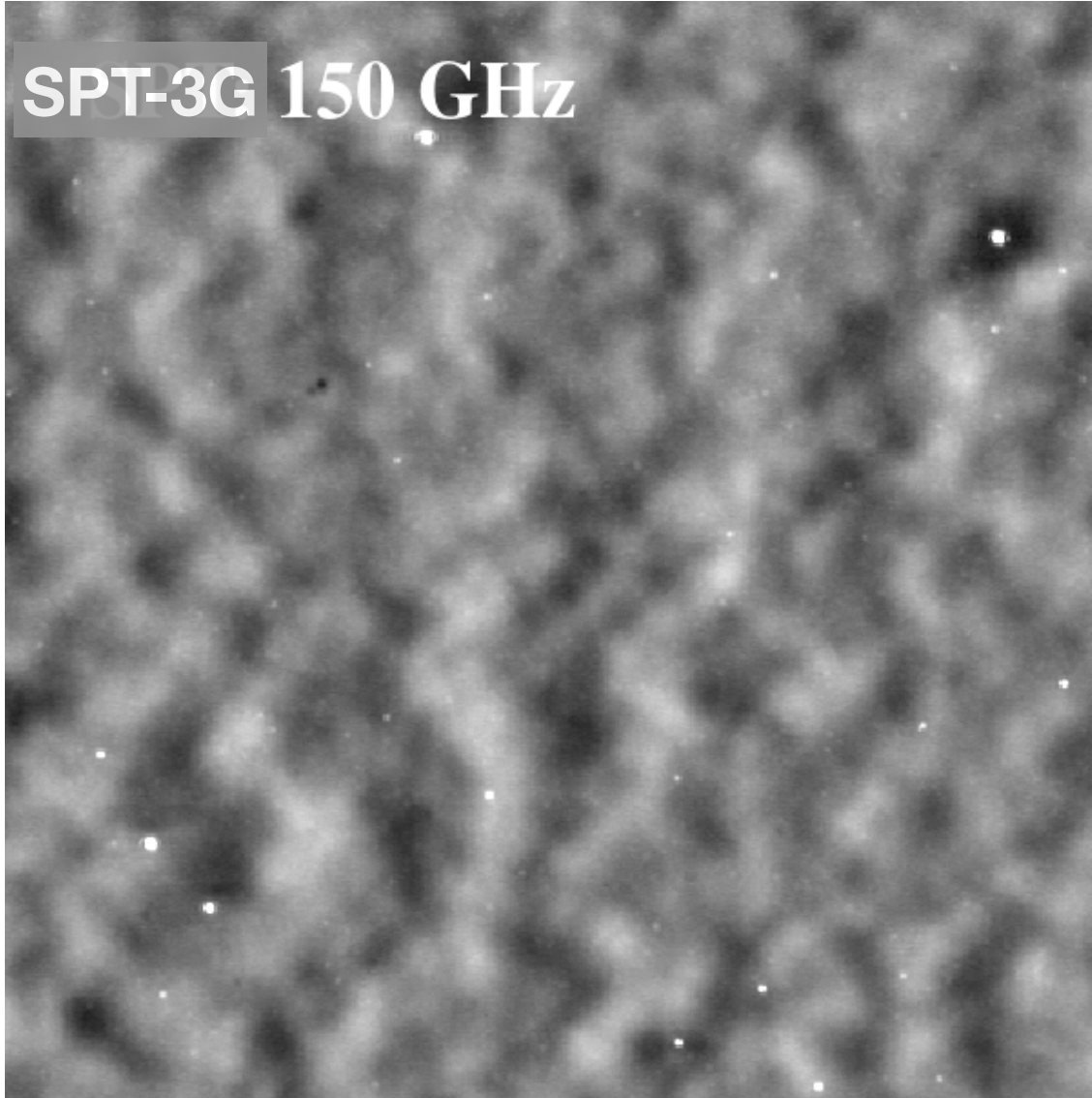


3 degrees

Srini Raghunathan

SPT-3G 150 GHz

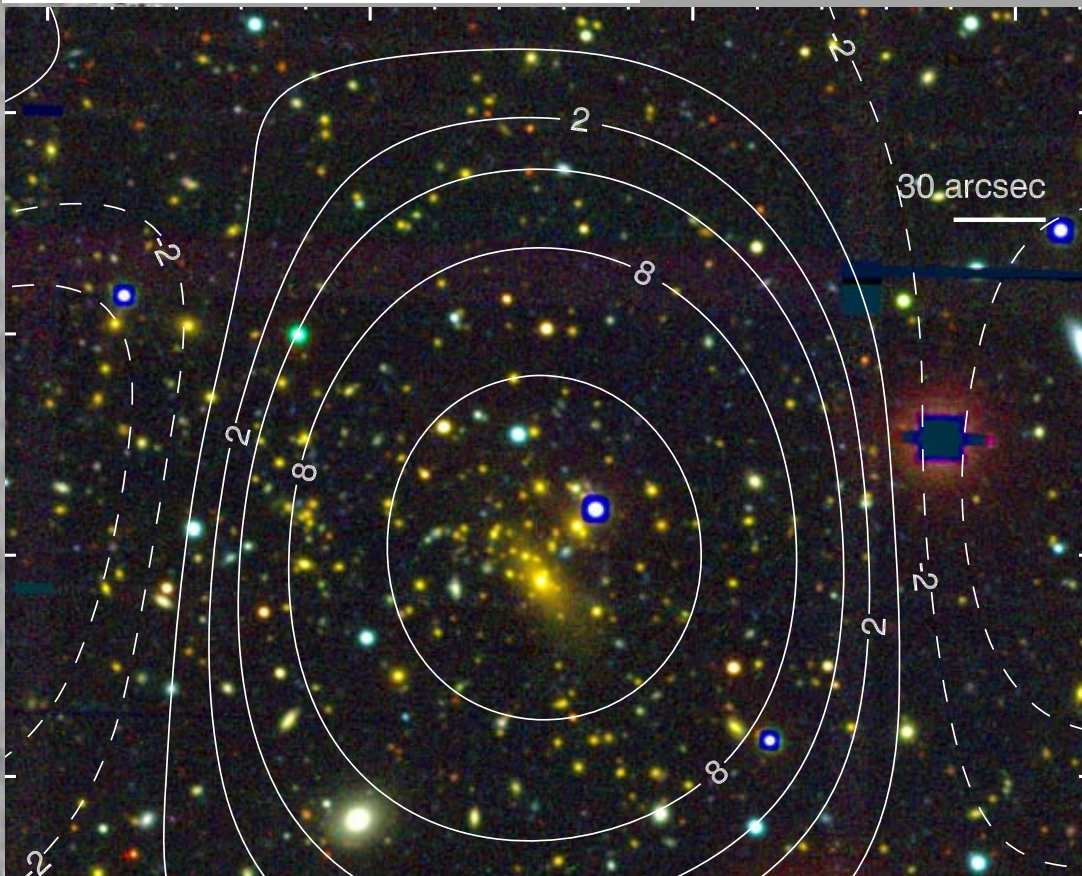
3 degrees



3 degrees

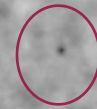
Srini Raghunathan

Image by Will High in recent paper by Williamson et al



One of the heaviest objects in the universe
 $> 10^{15}$ solar masses

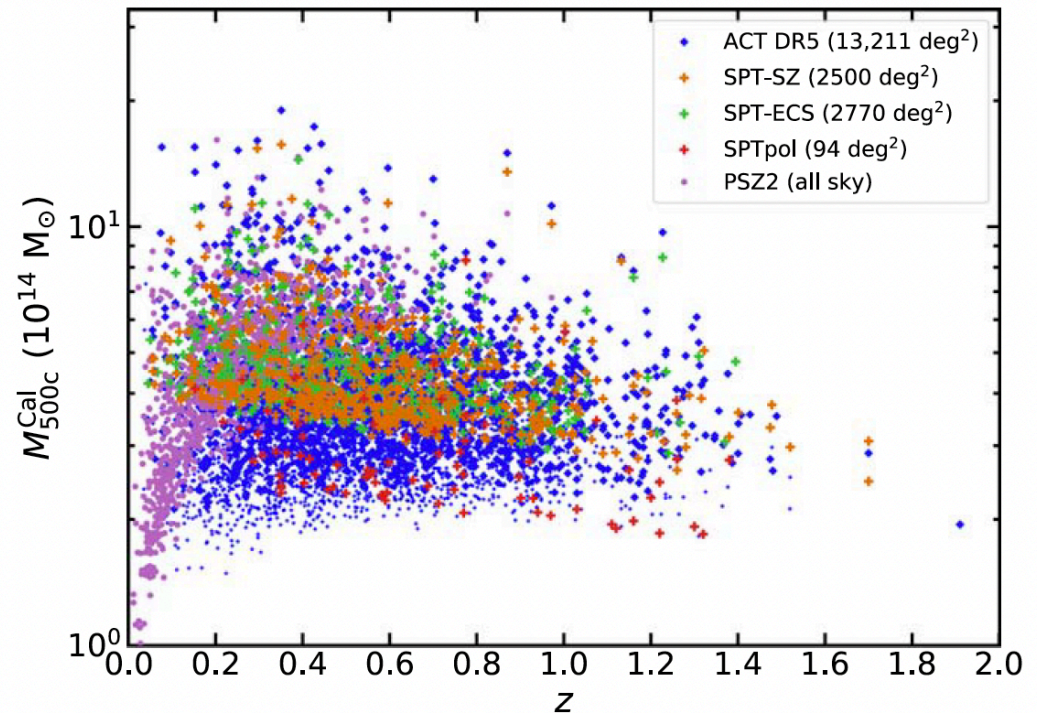
1 degree



patch of
isolated cosmic
fog

tSZ-selected Galaxy Clusters

- now many thousands of galaxy clusters have been discovered by their CMB signatures

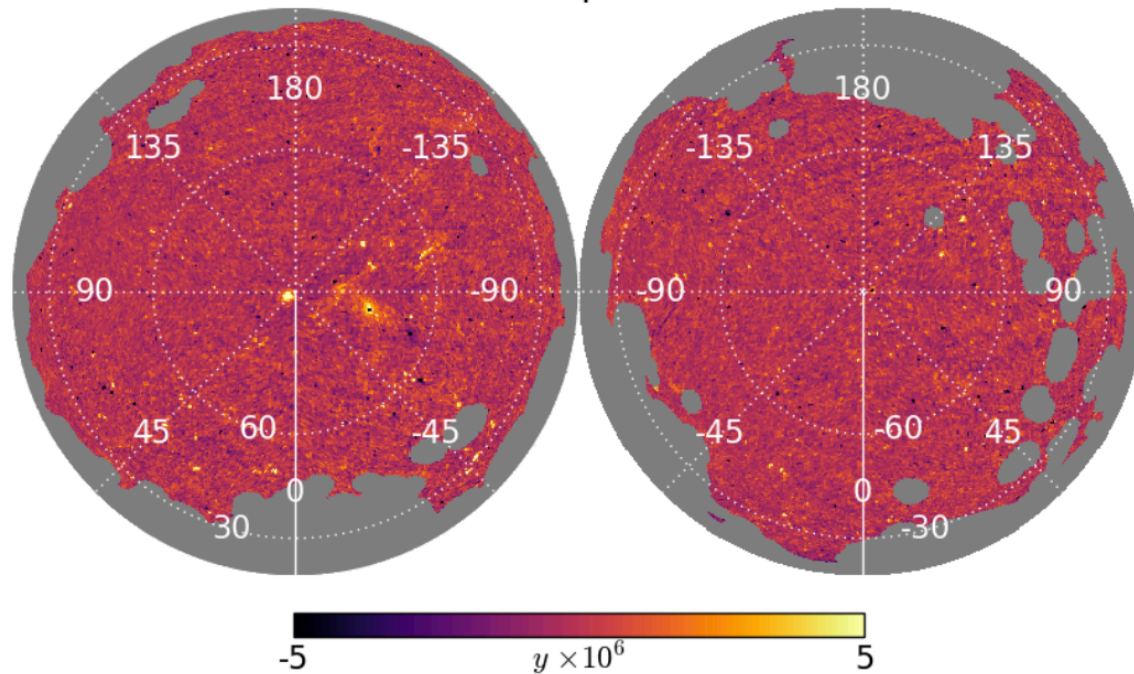


Hilton et al 2021

Compton y maps

Tanimura et al.

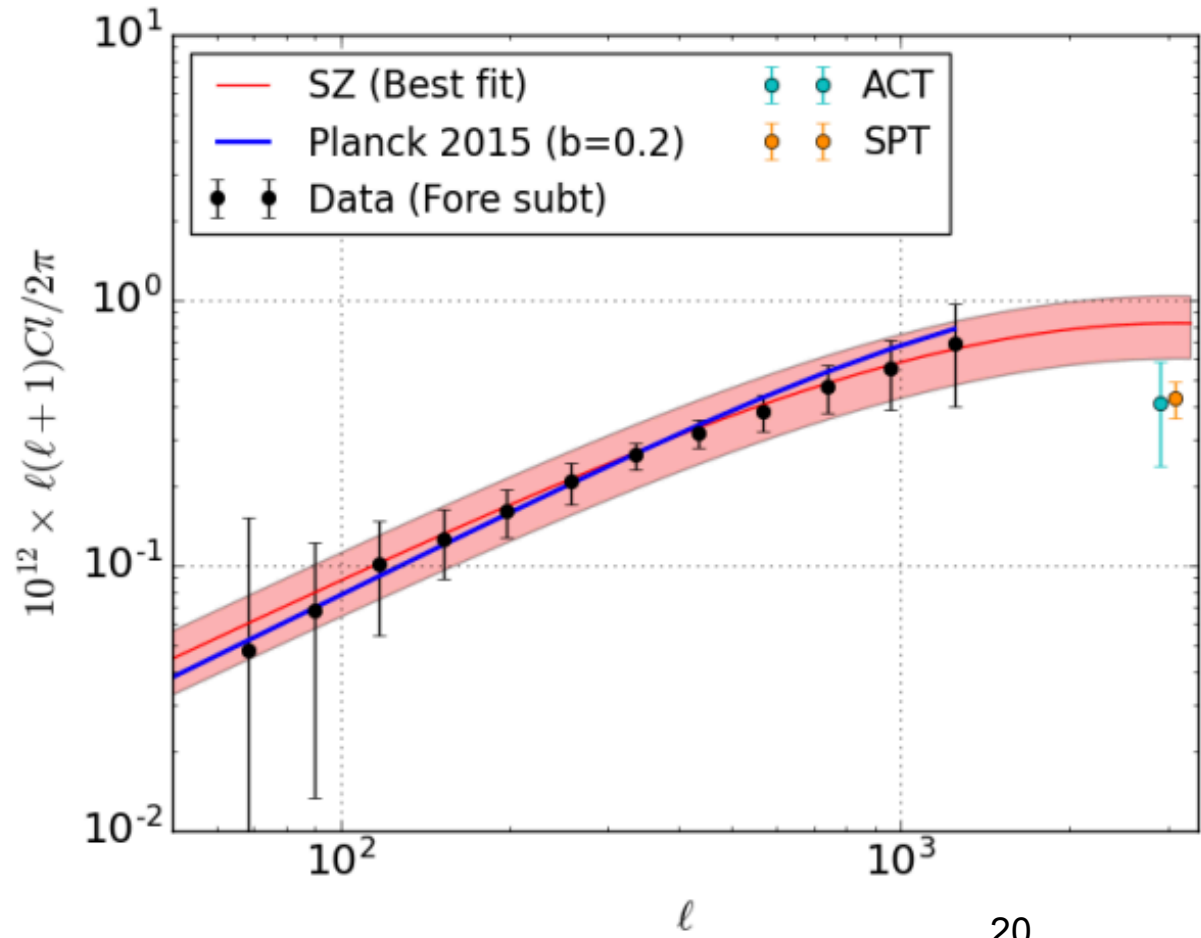
tSZ map 2020



Compton y power spectrum

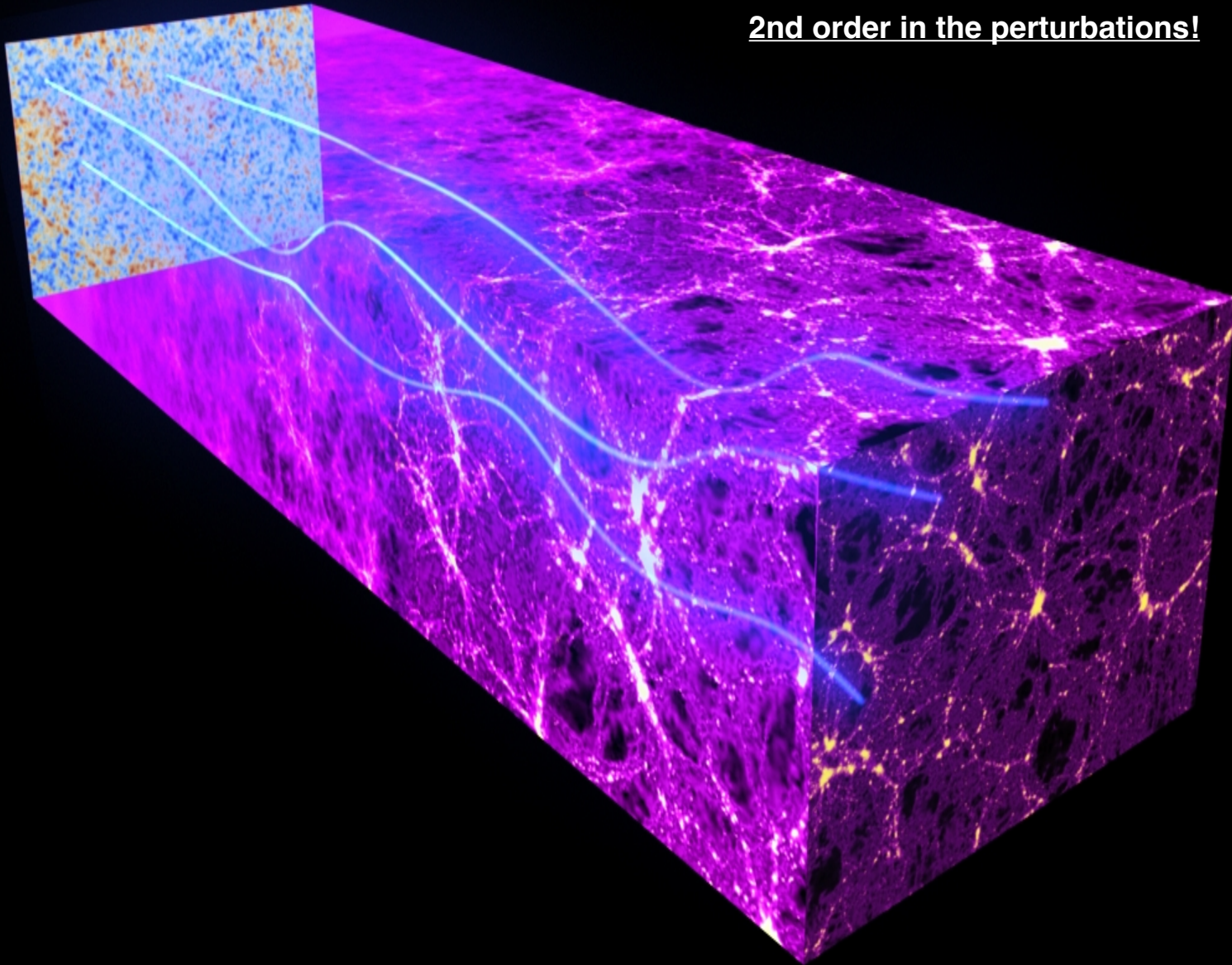
Hint that maybe tSZ power is low at high ℓ

Almost entirely just l-halo term



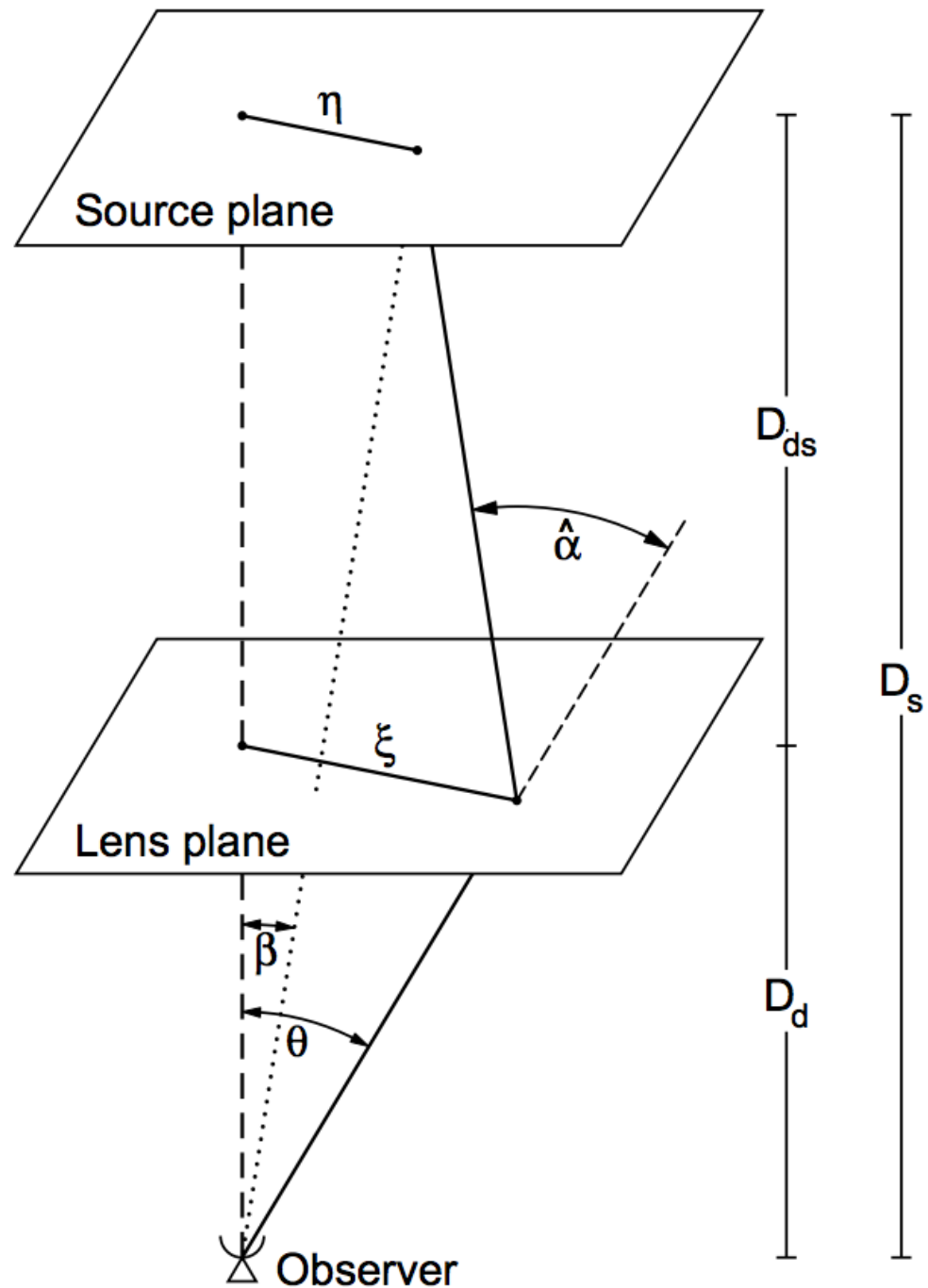
lensing of primordial fluctuations by intervening fluctuations

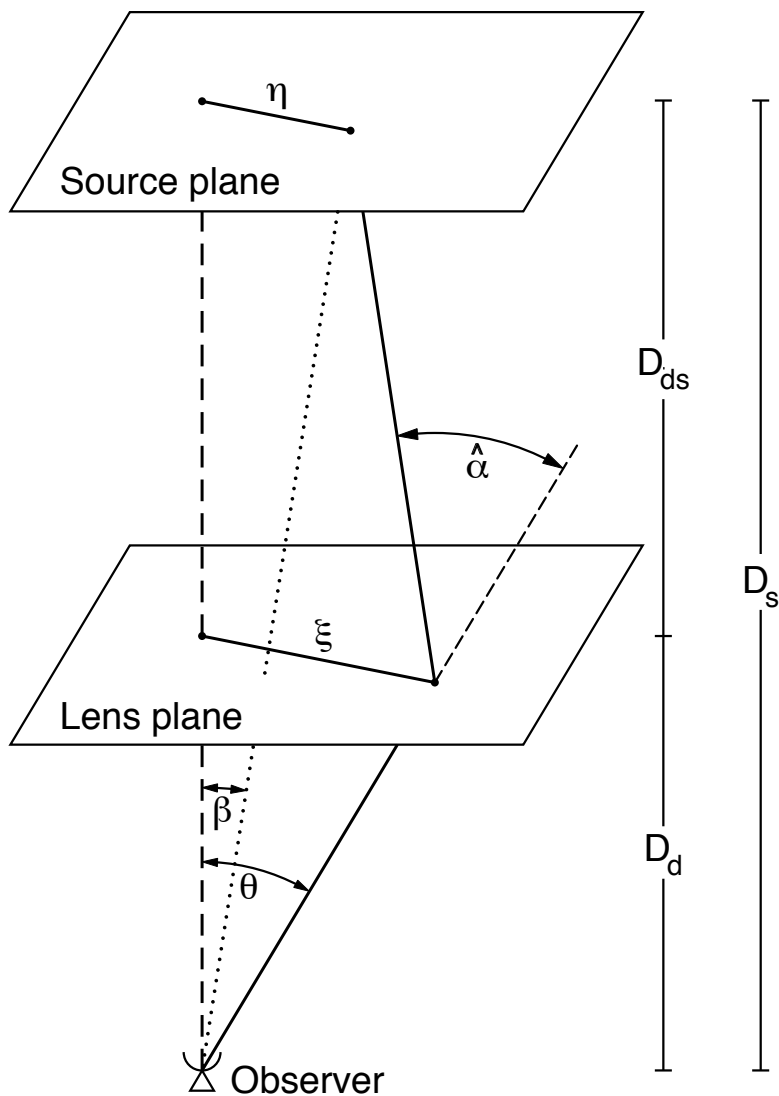
2nd order in the perturbations!



Gravitational lensing

$$\hat{\alpha} = \frac{4GM}{c^2 \xi}$$





source
position

image
position

deflection
angle

$$\vec{\beta} = \vec{\theta} - \frac{D_{ds}}{D_s} \hat{\alpha}(D_d \vec{\theta}) \equiv \vec{\theta} - \vec{\alpha}(\vec{\theta})$$

$$\vec{\alpha}(\vec{\theta}) = \frac{1}{\pi} \int_{\mathbb{R}^2} d^2\theta' \kappa(\vec{\theta}') \frac{\vec{\theta} - \vec{\theta}'}{|\vec{\theta} - \vec{\theta}'|^2}$$

$$\psi(\vec{\theta}) = \frac{1}{\pi} \int_{\mathbb{R}^2} d^2\theta' \kappa(\vec{\theta}') \ln |\vec{\theta} - \vec{\theta}'|$$

lensing
potential

convergence

$$\mathcal{A}(\vec{\theta}) = \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \left(\delta_{ij} - \frac{\partial^2 \psi(\vec{\theta})}{\partial \theta_i \partial \theta_j} \right) = \begin{pmatrix} 1 - \kappa - \gamma_1 & -\gamma_2 \\ -\gamma_2 & 1 - \kappa + \gamma_1 \end{pmatrix},$$

where we have introduced the components of the shear $\gamma \equiv \gamma_1 + i\gamma_2 = |\gamma|e^{2i\varphi}$

$$\gamma_1 = \frac{1}{2}(\psi_{,11} - \psi_{,22}), \quad \gamma_2 = \psi_{,12},$$

$$\mathcal{A} = (1 - \kappa) \begin{pmatrix} 1 - g_1 & -g_2 \\ -g_2 & 1 + g_1 \end{pmatrix}$$

distortion has overall magnification

image gets bigger (or smaller),
not **brighter** (dimmer)

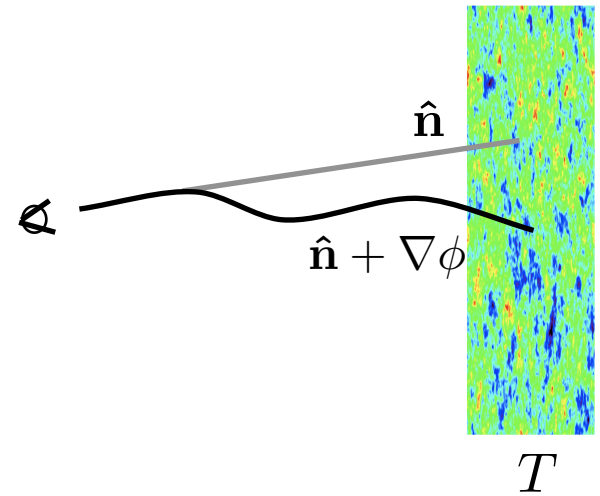
$$g(\vec{\theta}) \equiv \frac{\gamma(\vec{\theta})}{1 - \kappa(\vec{\theta})}$$

reduced shear

CMB Lensing

Photons get shifted

$$T^L(\hat{\mathbf{n}}) = T^U(\hat{\mathbf{n}} + \nabla\phi(\hat{\mathbf{n}}))$$



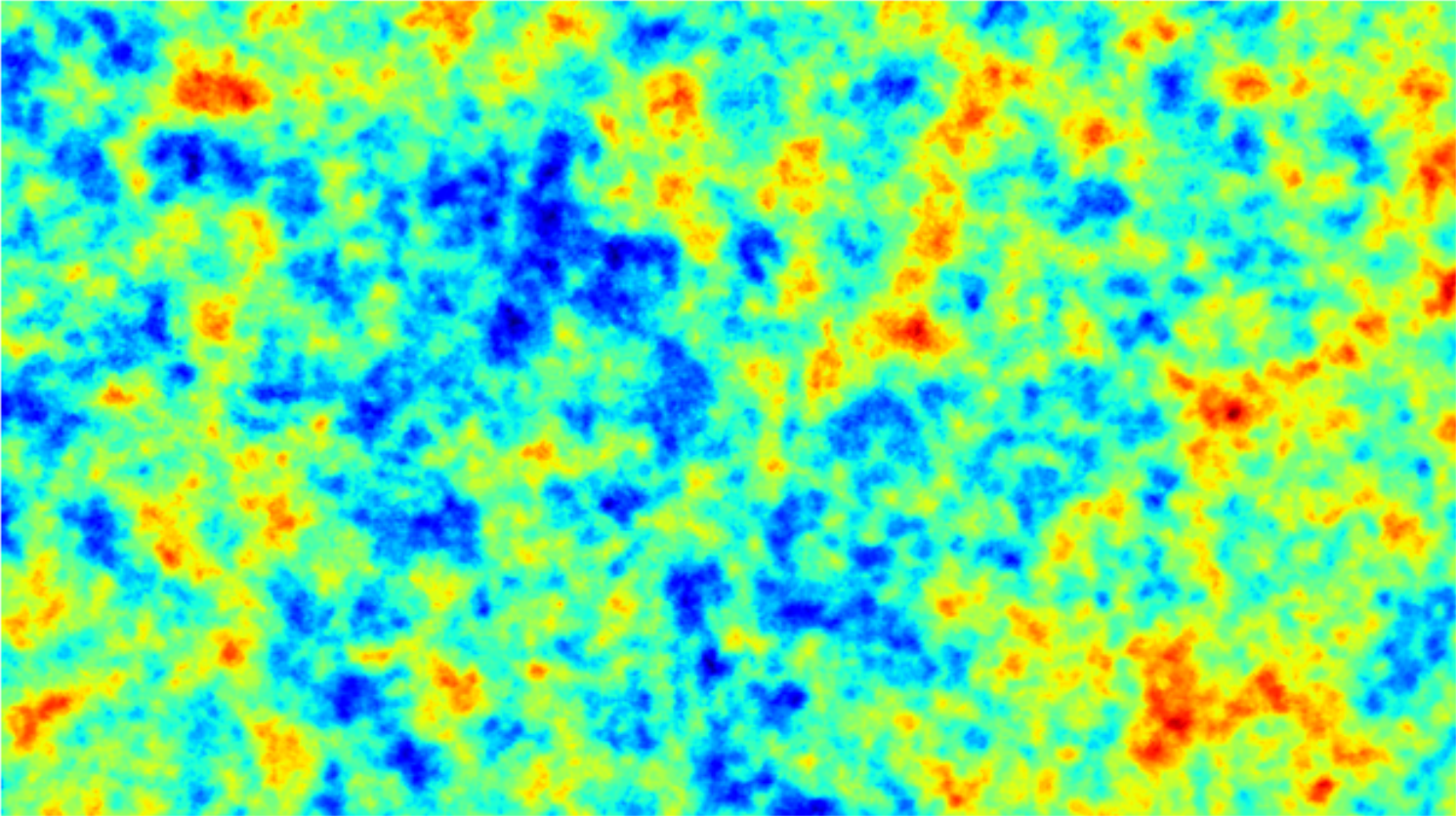
In WL limit, add many deflections along line of sight

$$\nabla\phi(\hat{\mathbf{n}}) = -2 \int_0^{\chi_*} d\chi \frac{\chi_* - \chi}{\chi_*\chi} \nabla_{\perp} \Phi(\chi\hat{\mathbf{n}}, \chi)$$

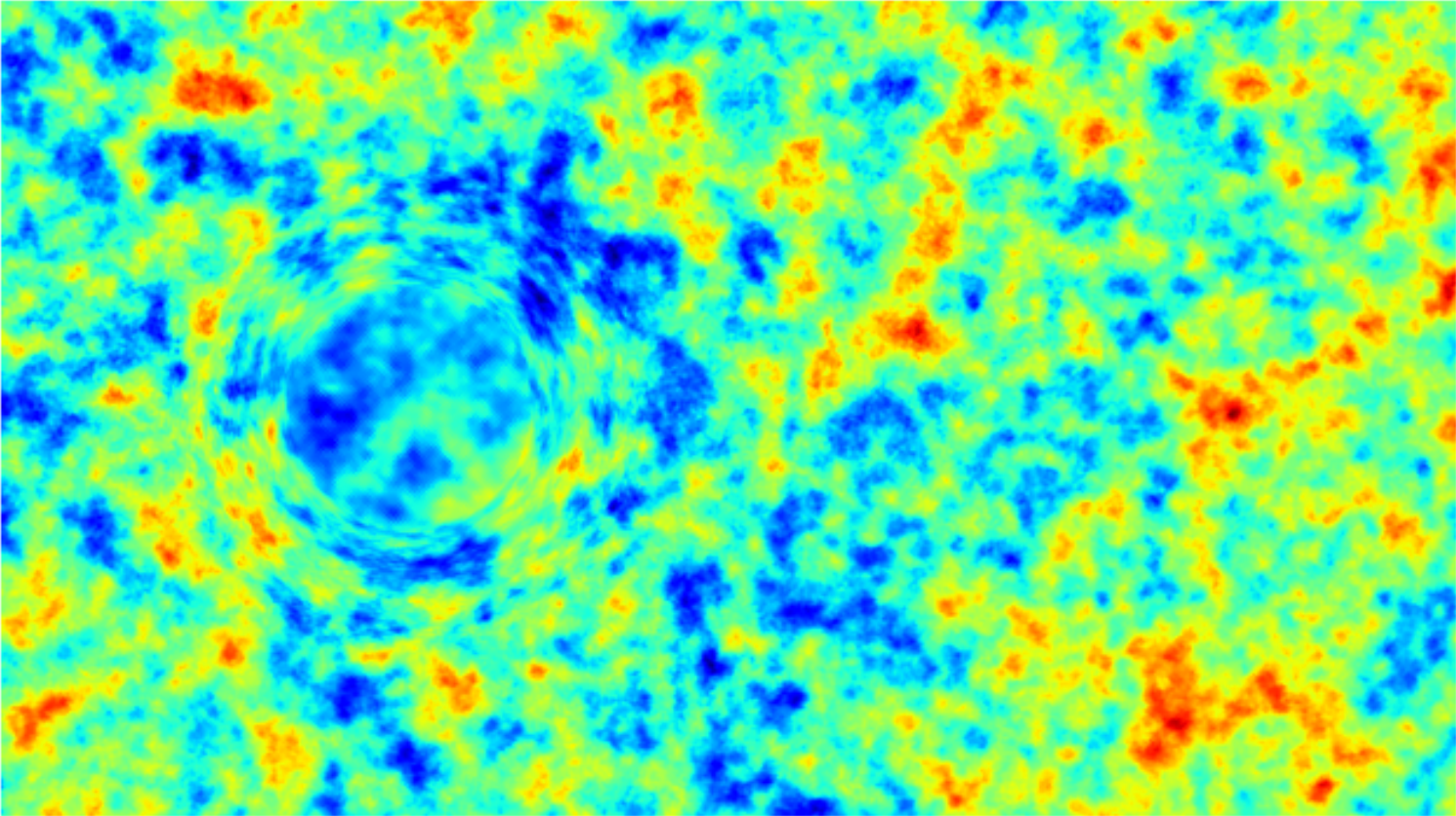
Broad kernel, peaks at $z \sim 2$

- CMB is a unique source for lensing
 - Gaussian, with well-understood power spectrum (contains all info)
 - At redshift which is (a) unique, (b) known, and (c) highest

patch of sky (the North pole) as seen by Planck (17x10 degrees)

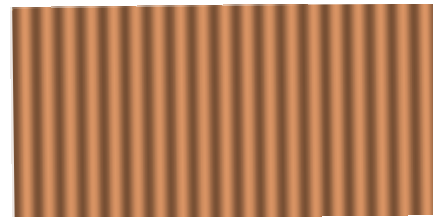
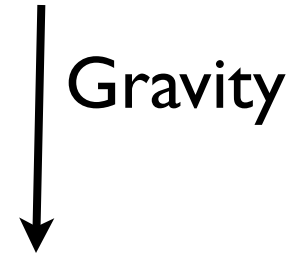
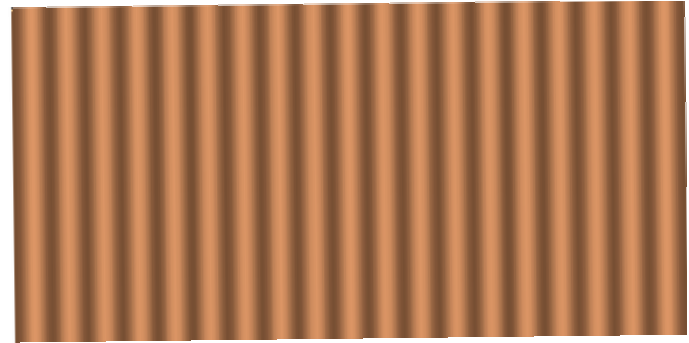


SIMULATED lensing effect (20x larger than typical)



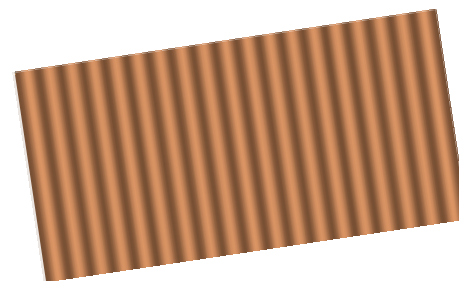
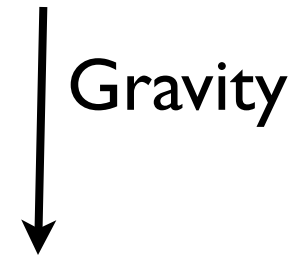
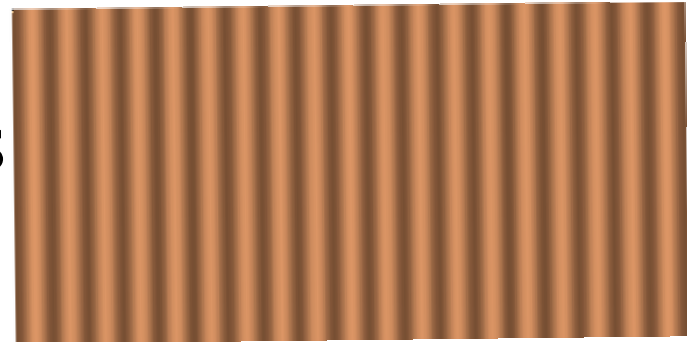
Lensing simplified

- gravitational potentials distort images by stretching, squeezing, shearing

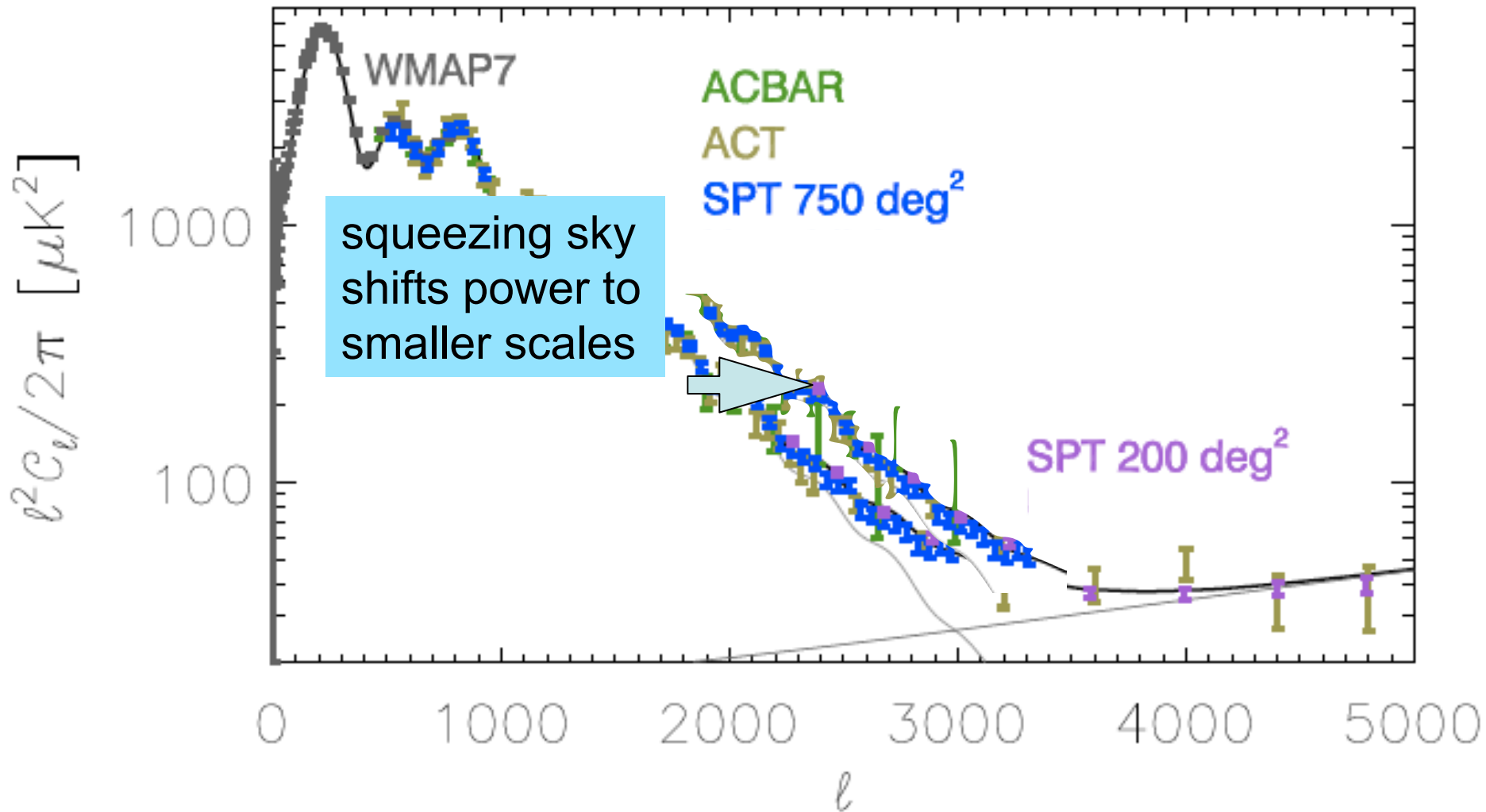


Lensing simplified

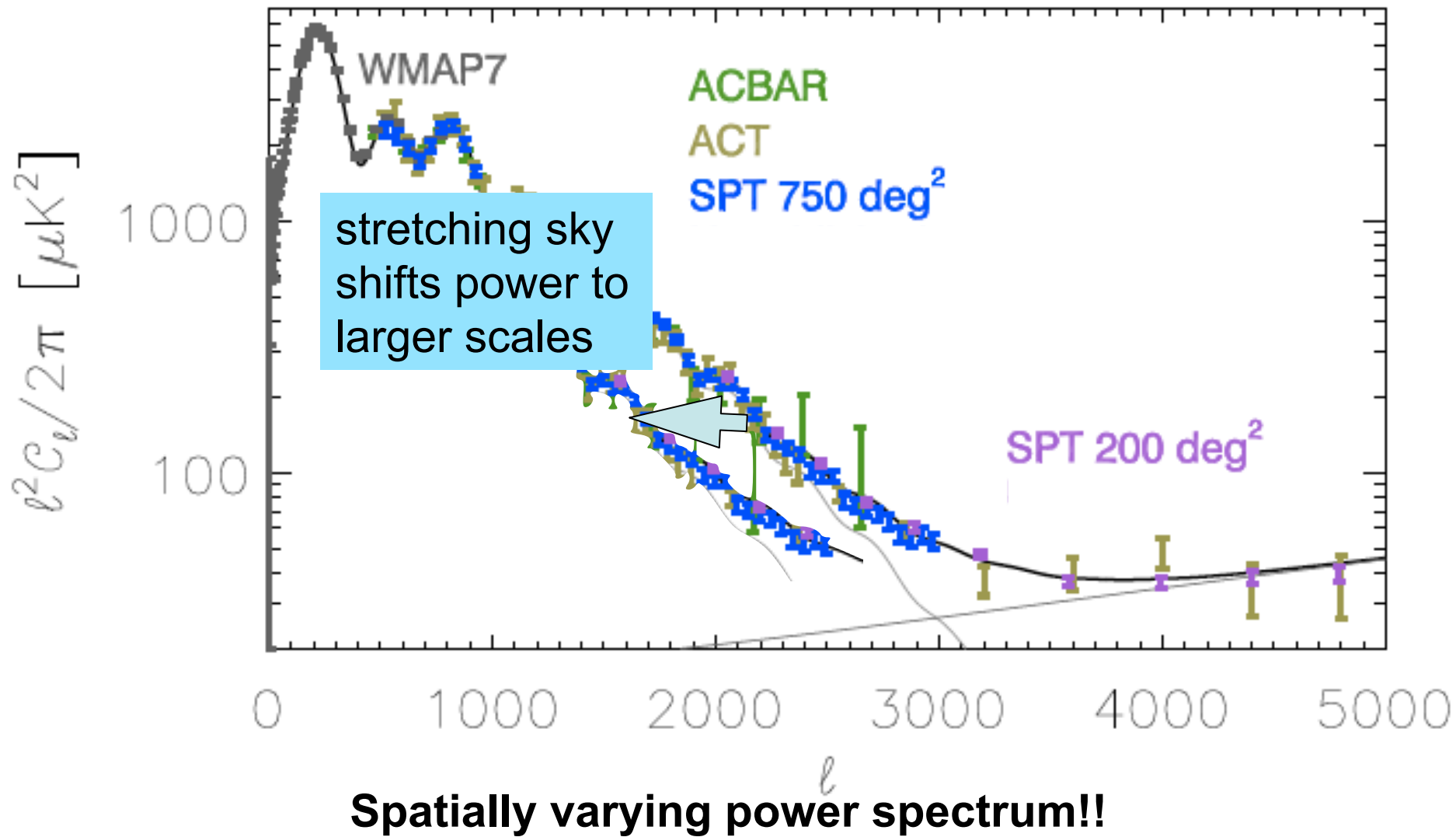
- where gravity stretches, gradients become smaller
- where gravity compresses, gradients are larger
- shear changes ***direction***



CMB Power Spectrum

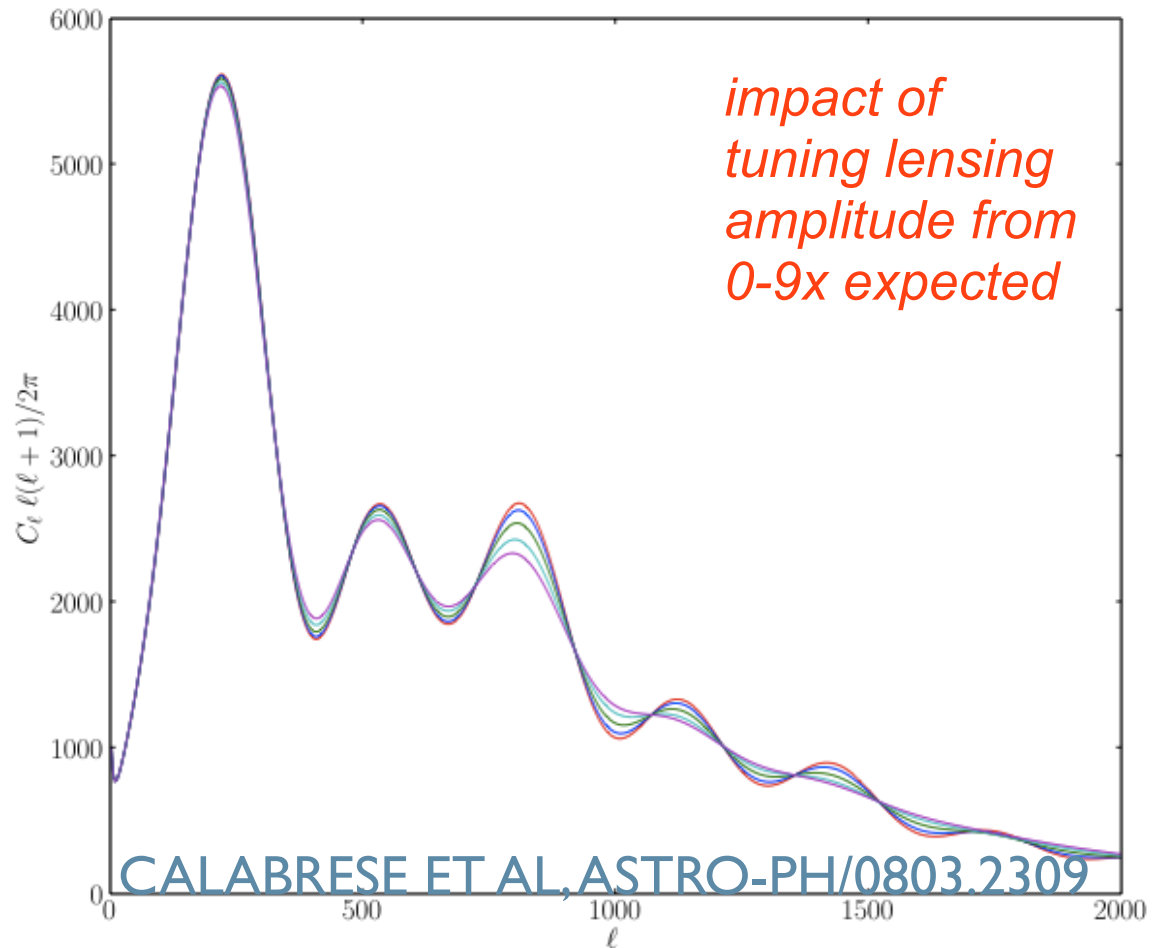


CMB Power Spectrum



Effect on CMB Power Spectrum

- mixing of power leads to smoothing of acoustic peaks
- small effect but data is really good



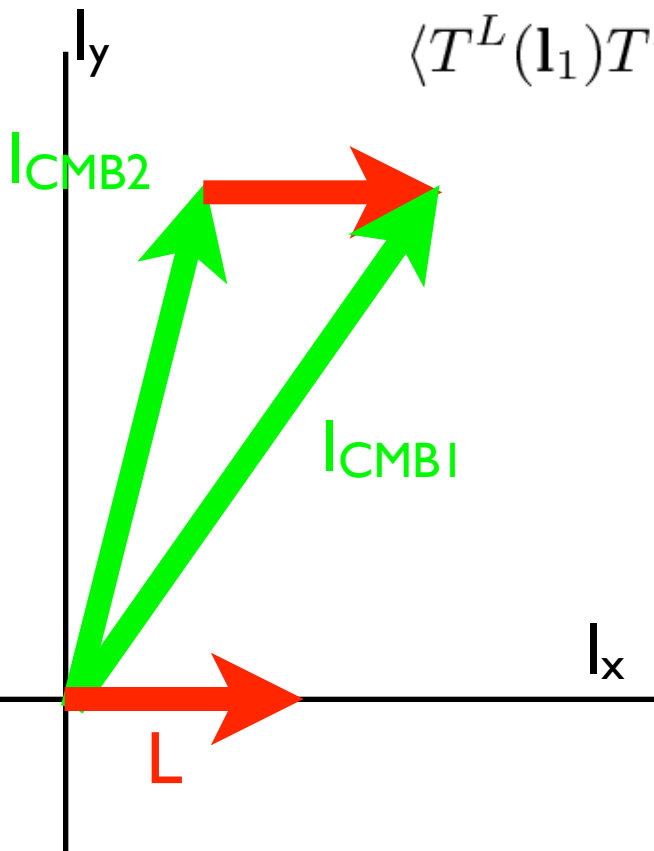
Mode Coupling from Lensing

$$\begin{aligned} T^L(\hat{\mathbf{n}}) &= T^U(\hat{\mathbf{n}} + \nabla\phi(\hat{\mathbf{n}})) \\ &= T^U(\hat{\mathbf{n}}) + \nabla T^U(\hat{\mathbf{n}}) \cdot \nabla\phi(\hat{\mathbf{n}}) + O(\phi^2), \end{aligned}$$

- Non-gaussian mode coupling for $\mathbf{l}_1 \neq -\mathbf{l}_2$:

$$\langle T^L(\mathbf{l}_1)T^L(\mathbf{l}_2) \rangle = \mathbf{L} \cdot (\mathbf{l}_1 C_{l_1}^T + \mathbf{l}_2 C_{l_2}^T) \phi(\mathbf{L}) + O(\phi^2)$$

$$\mathbf{L} = \mathbf{l}_1 + \mathbf{l}_2$$



- We extract ϕ by taking a suitable average over CMB multipoles separated by a distance L
- We use the standard Hu quadratic estimator.

E-modes and B-modes

$$Q(l) = [E(l) \cos(2\phi_l) - B(l) \sin(2\phi_l)]$$

$$U(l) = [E(l) \sin(2\phi_l) + B(l) \cos(2\phi_l)].$$

- E/B is a different way to express polarization field
- easy to understand in flat-sky limit (i.e. Fourier modes)

E-modes/B-modes

- E-modes vary spatially parallel or perpendicular to polarization direction
- B-modes vary spatially at 45 degrees
- CMB
 - scalar perturbations only generate *only* E
- ***Lensing of CMB is much more obvious in polarization!***

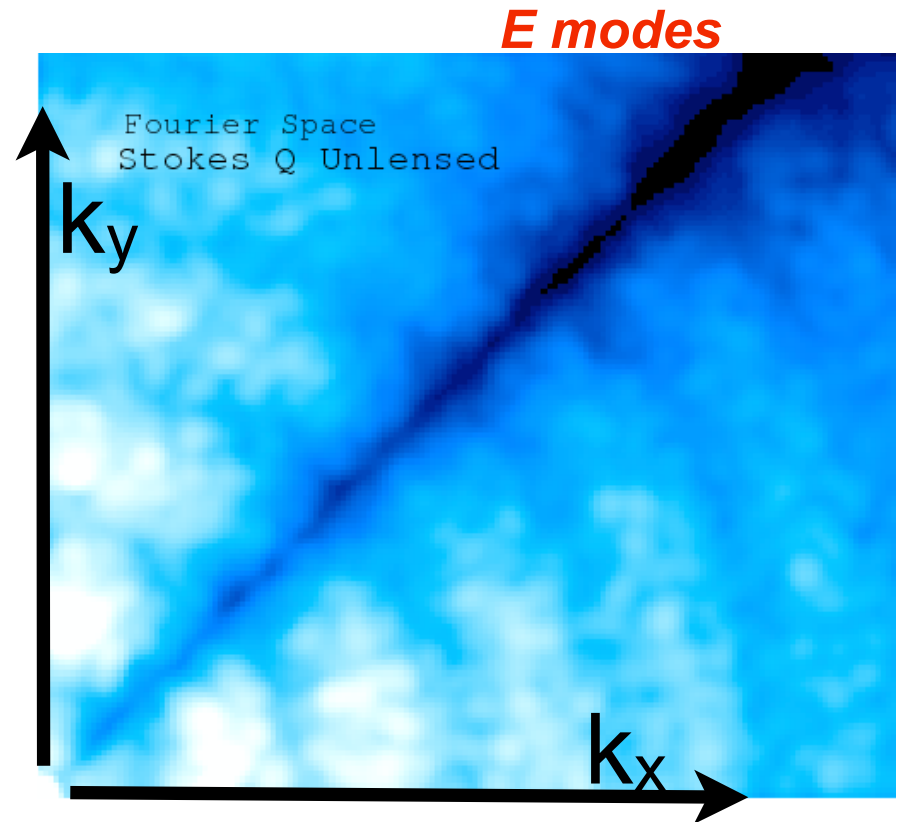
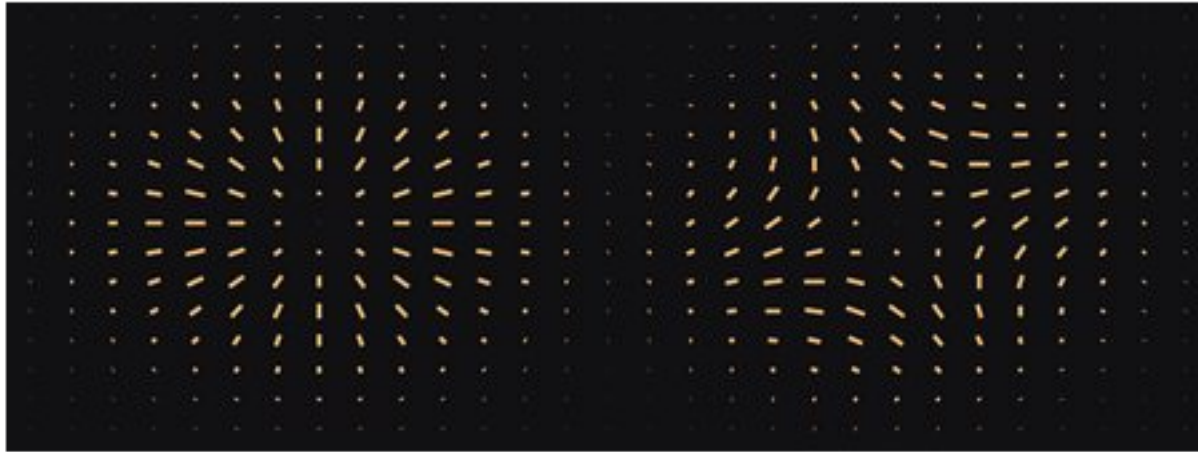


Image of positive k_x /positive k_y Fourier transform of a 10x10 deg chunk of Stokes Q CMB map [simulated; nothing clever done to it]

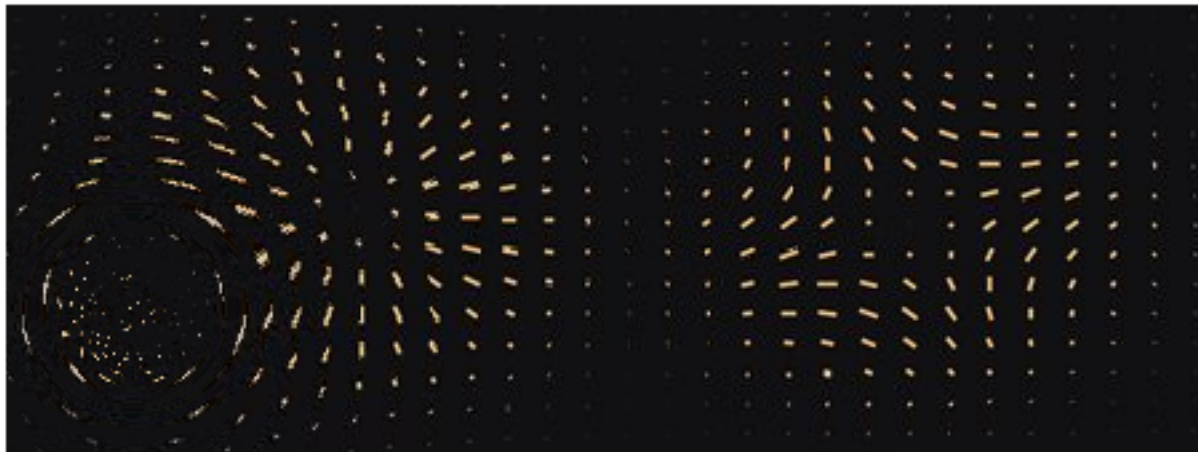
B Modes from E Modes

Before: pure E mode (left) and pure B mode (right)



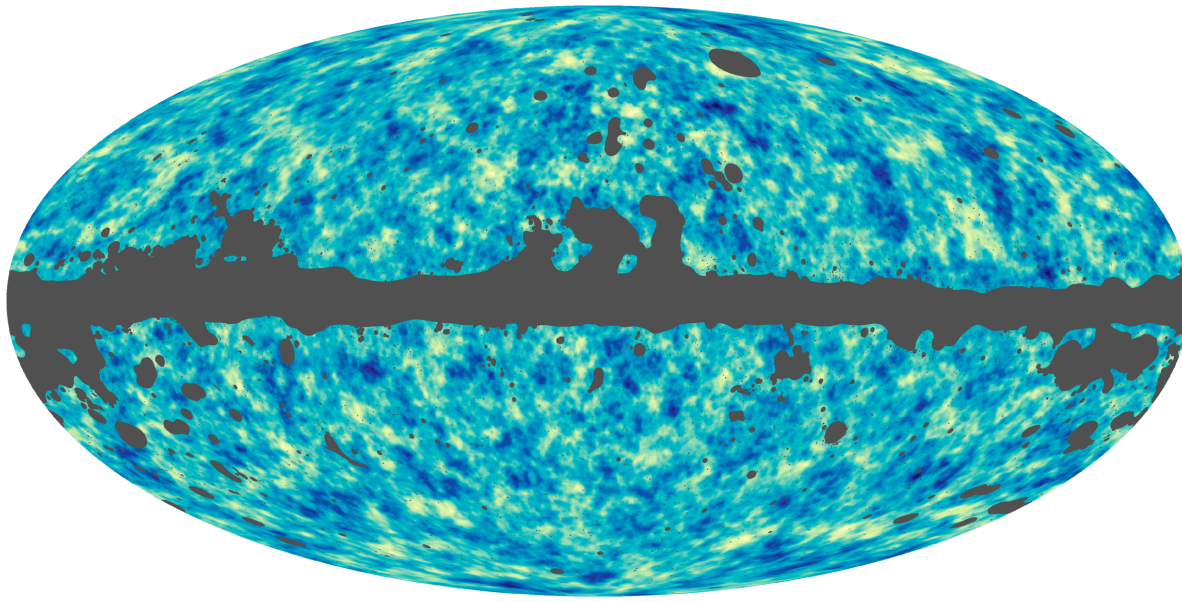
From B-pol.org

After: large point mass lenses image

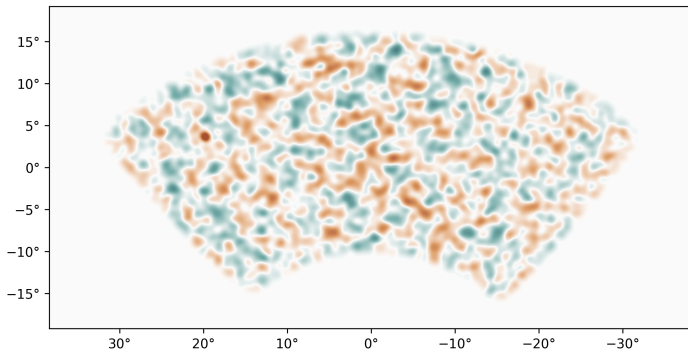


Lensing done with “Lens an astrophysicist”

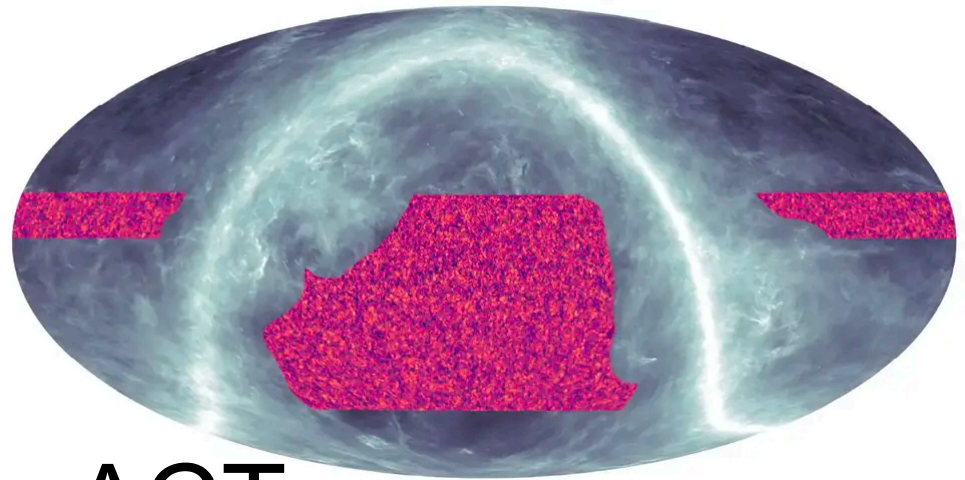
<http://theory2.phys.cwru.edu/~pete/GravitationalLens/>



Planck
(~all-sky)



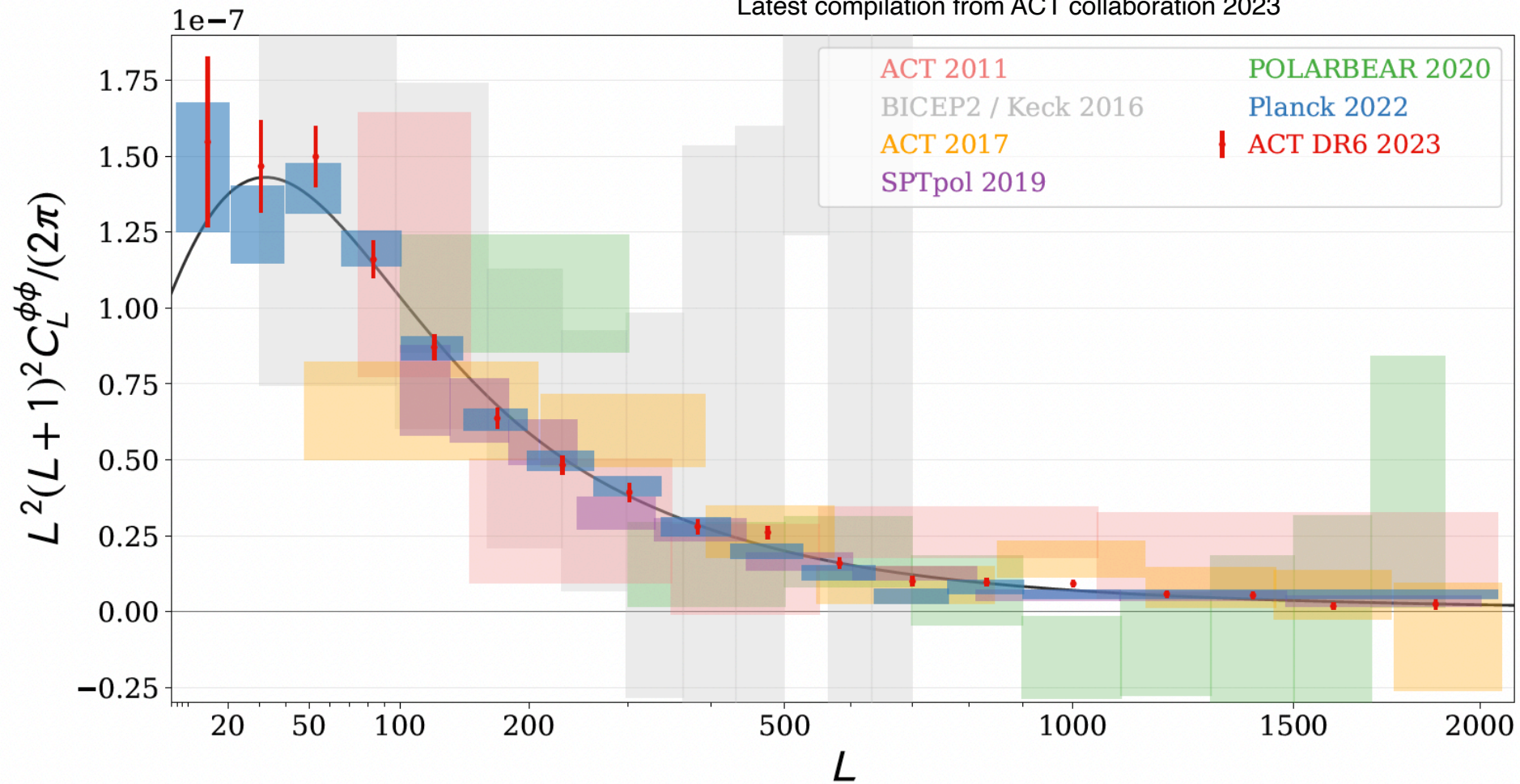
SPT-3G
(1500 square degrees)



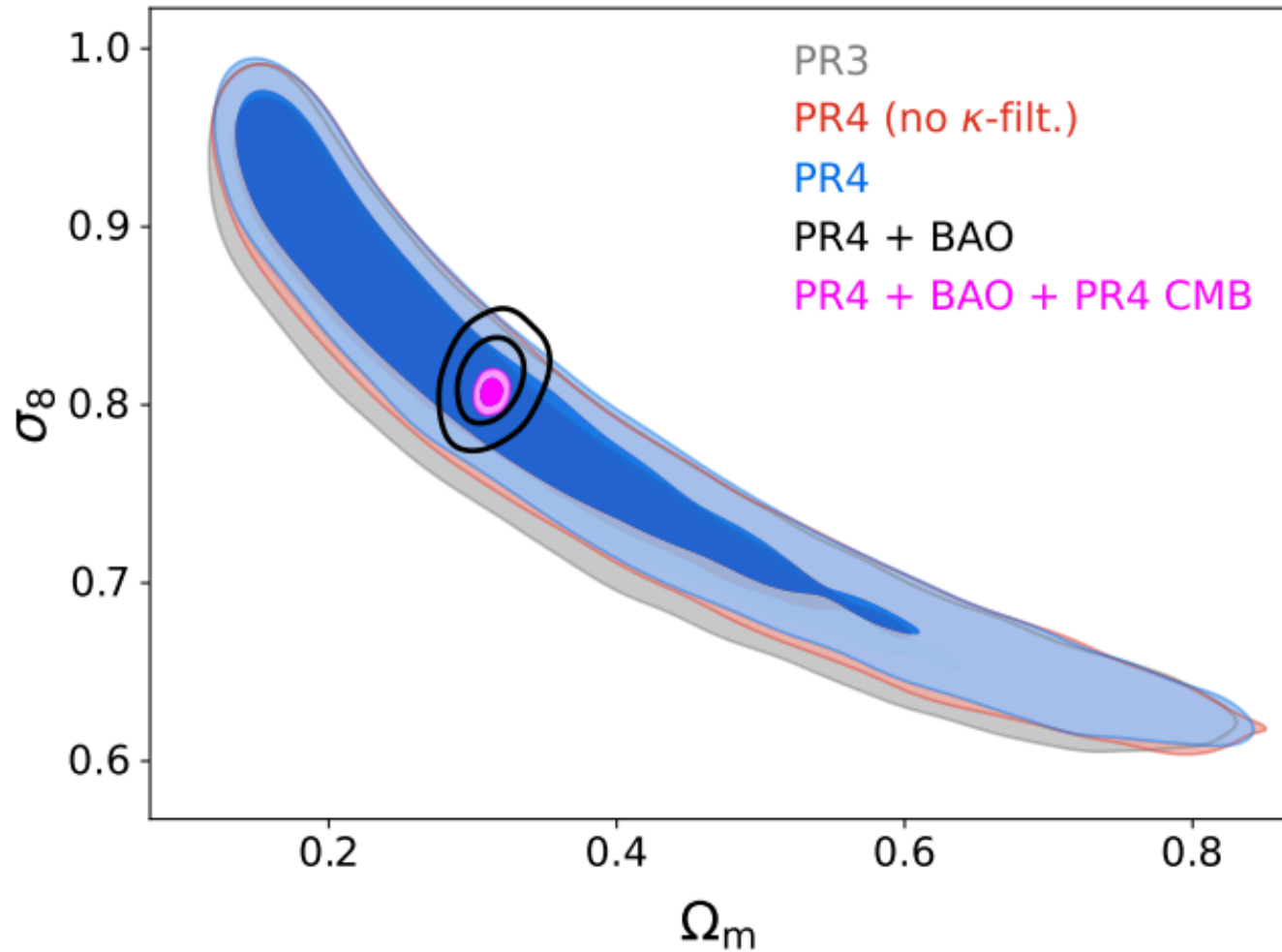
ACT
(9400 square degrees)

CMB Lensing Power Spectra

Latest compilation from ACT collaboration 2023



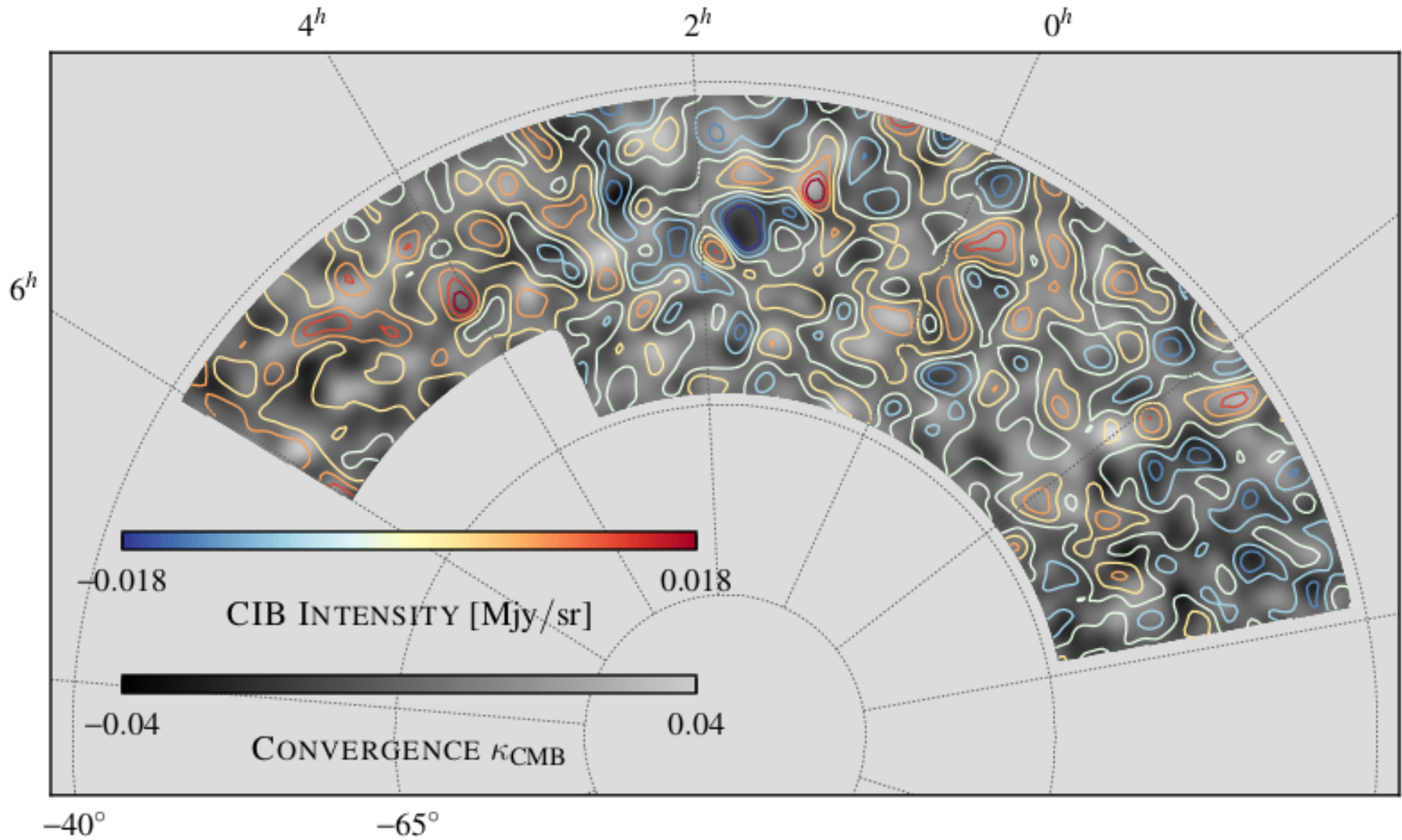
Cosmological constraints on structure formation



Planck: Carron 2022

see also Madhavacheril 2023

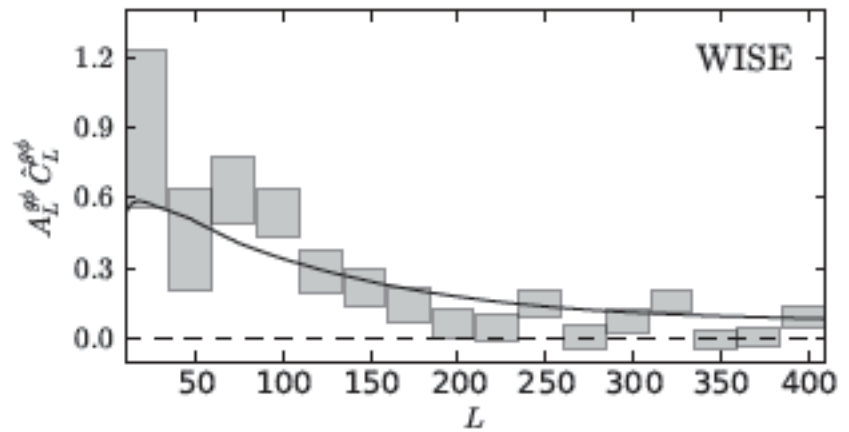
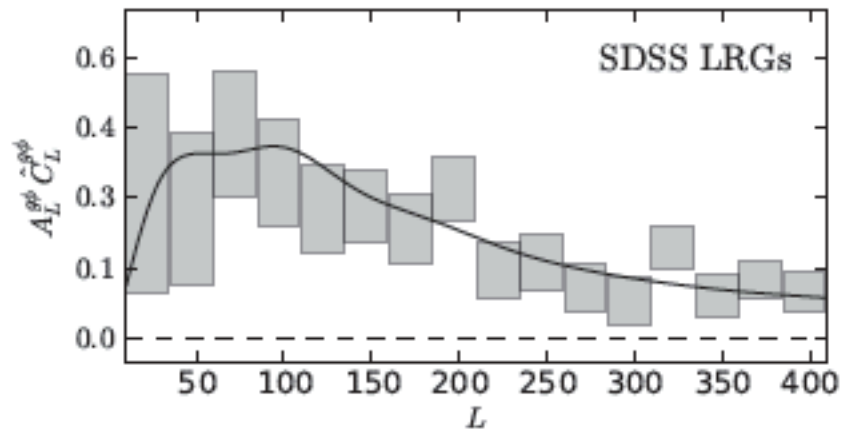
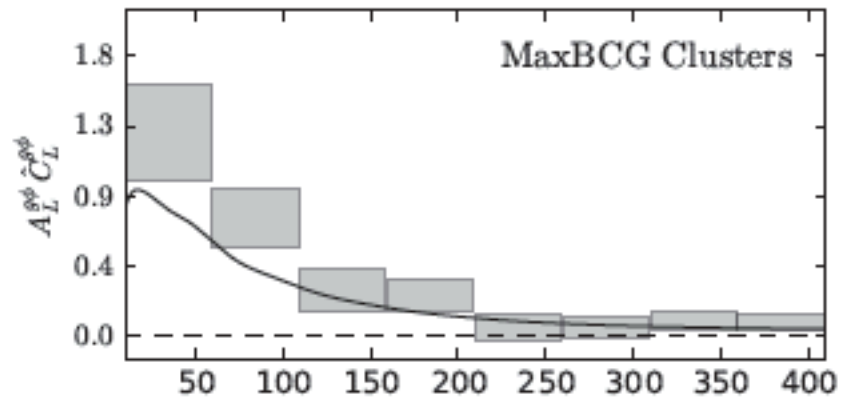
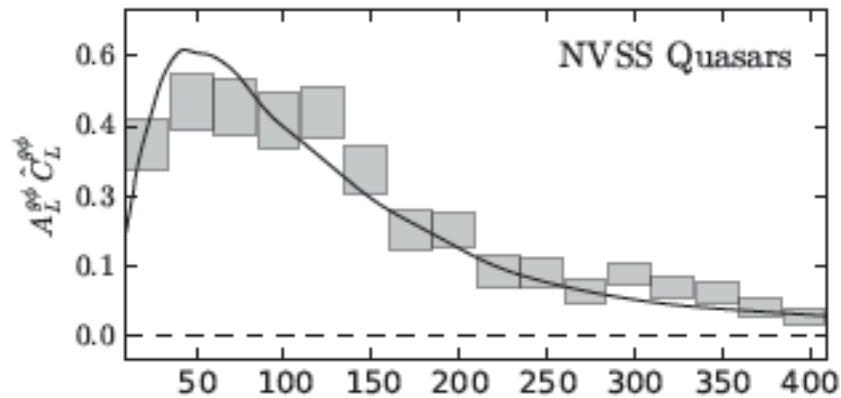
CMB-LSS cross-correlation: CIB



CIB map from Planck GNLIC 545 GHz

Omori, Chown, Simard, KTS, et. al (arXv:1705.00743)

Planck X Galaxies, etc.



Angular Clustering

Angular power spectrum of power spectrum between two maps X & Y (could be same map!)

$$C_\ell^{XY} = \frac{2}{\pi} \int_0^\infty d\chi_1 d\chi_2 W^X(\chi_1) W^Y(\chi_2) \int_0^\infty k^2 dk P_{XY}(k; z_1, z_2) j_\ell(k\chi_1) j_\ell(k\chi_2)$$

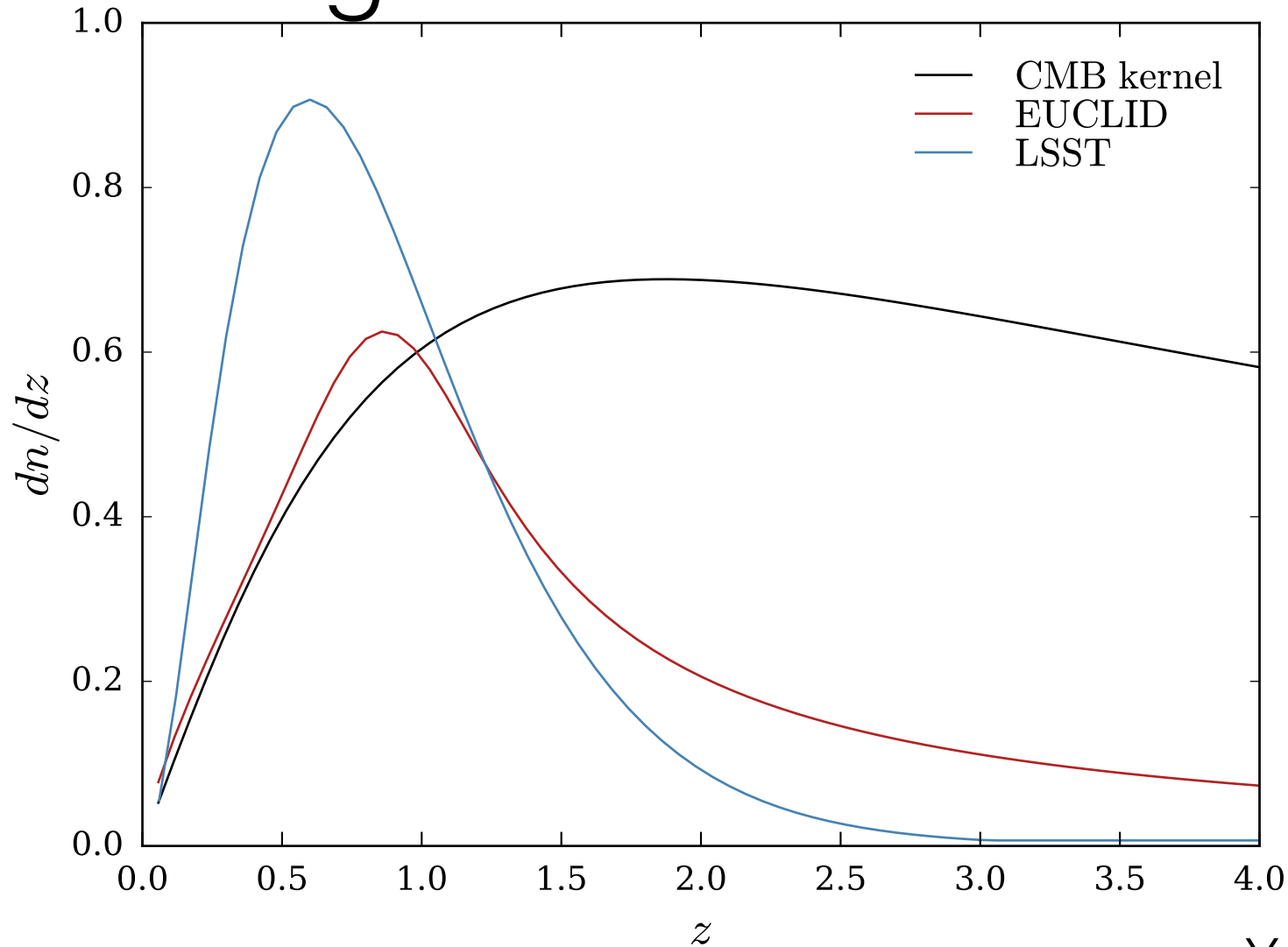
Limber approximation, which generally works pretty well except for really large scales

$$C_\ell^{XY} = \int d\chi \frac{W^X(\chi) W^Y(\chi)}{\chi^2} P_{XY} \left(k_\perp = \frac{\ell + 1/2}{\chi}, k_z = 0 \right)$$

weights for CMB lensing or some galaxy tracer

$$W^\kappa(\chi) = \frac{3}{2} (\Omega_m + \Omega_\nu) H_0^2 (1+z) \frac{\chi(\chi_\star - \chi)}{\chi_\star}, \quad W^g(\chi) = b(z) H(z) \frac{dN}{dz}$$

CMB lensing is sensitive to higher z sources



Angular Clustering

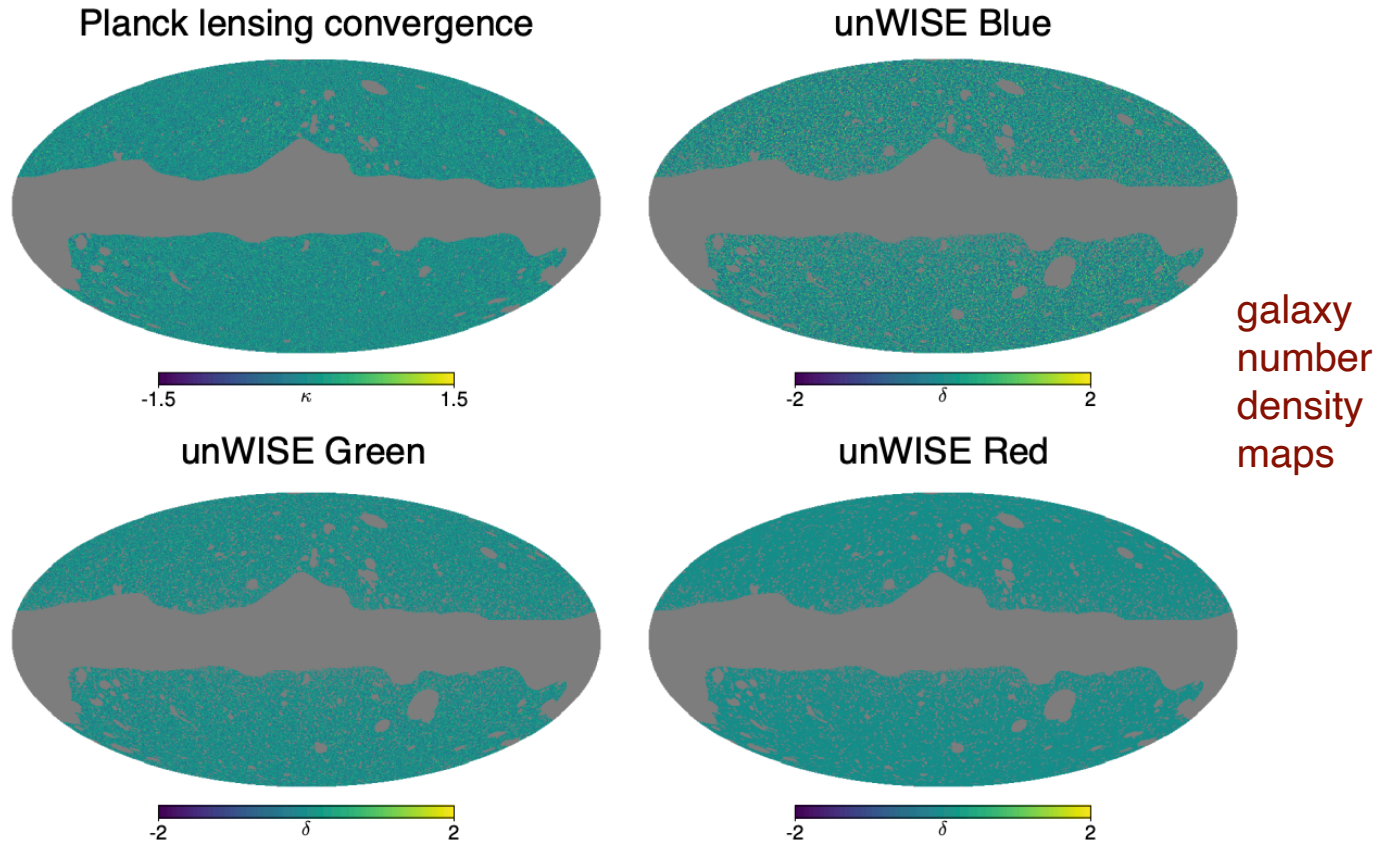
$$C_{\ell}^{\kappa g} = b^{\text{eff}} \int d\chi \frac{W^{\kappa}(\chi)}{\chi} H(z) \left[f(z) \frac{dN_p}{dz} \right] P(k\chi = \ell + 1/2)$$

$$C_{\ell}^{gg} = (b^{\text{eff}})^2 \int d\chi \frac{1}{\chi^2} H(z)^2 \left[f(z) \frac{dN_p}{dz} \right]^2 P(k\chi = \ell + 1/2)$$

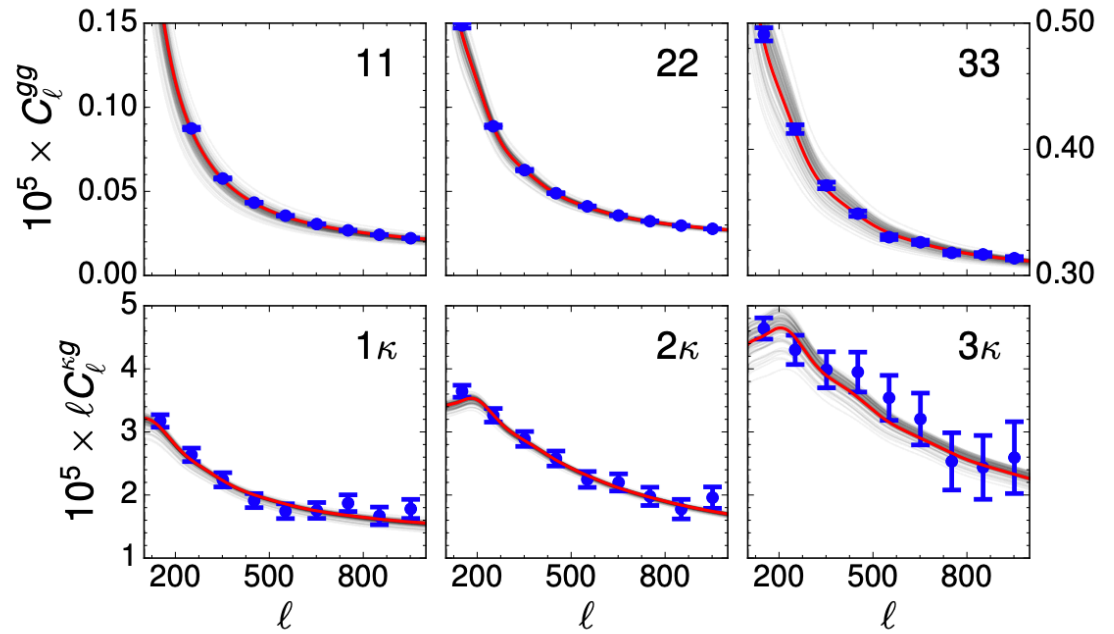
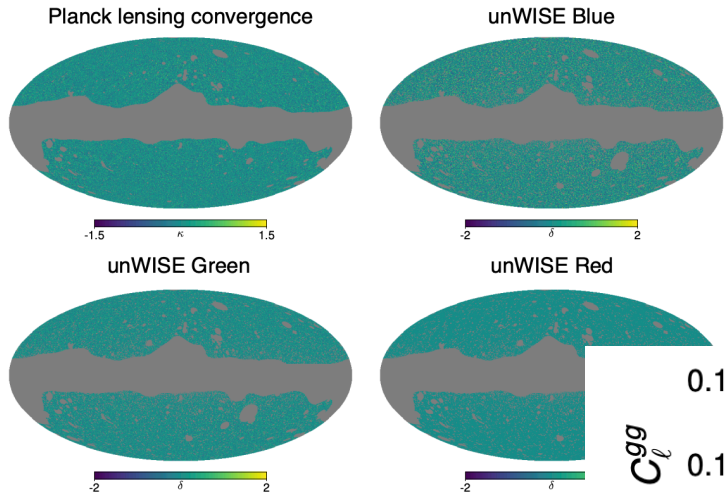
$$W^{\kappa}(\chi) = \frac{3}{2} (\Omega_m + \Omega_{\nu}) H_0^2 (1+z) \frac{\chi(\chi_{\star} - \chi)}{\chi_{\star}} \quad , \quad W^g(\chi) = b(z) H(z) \frac{dN}{dz}$$

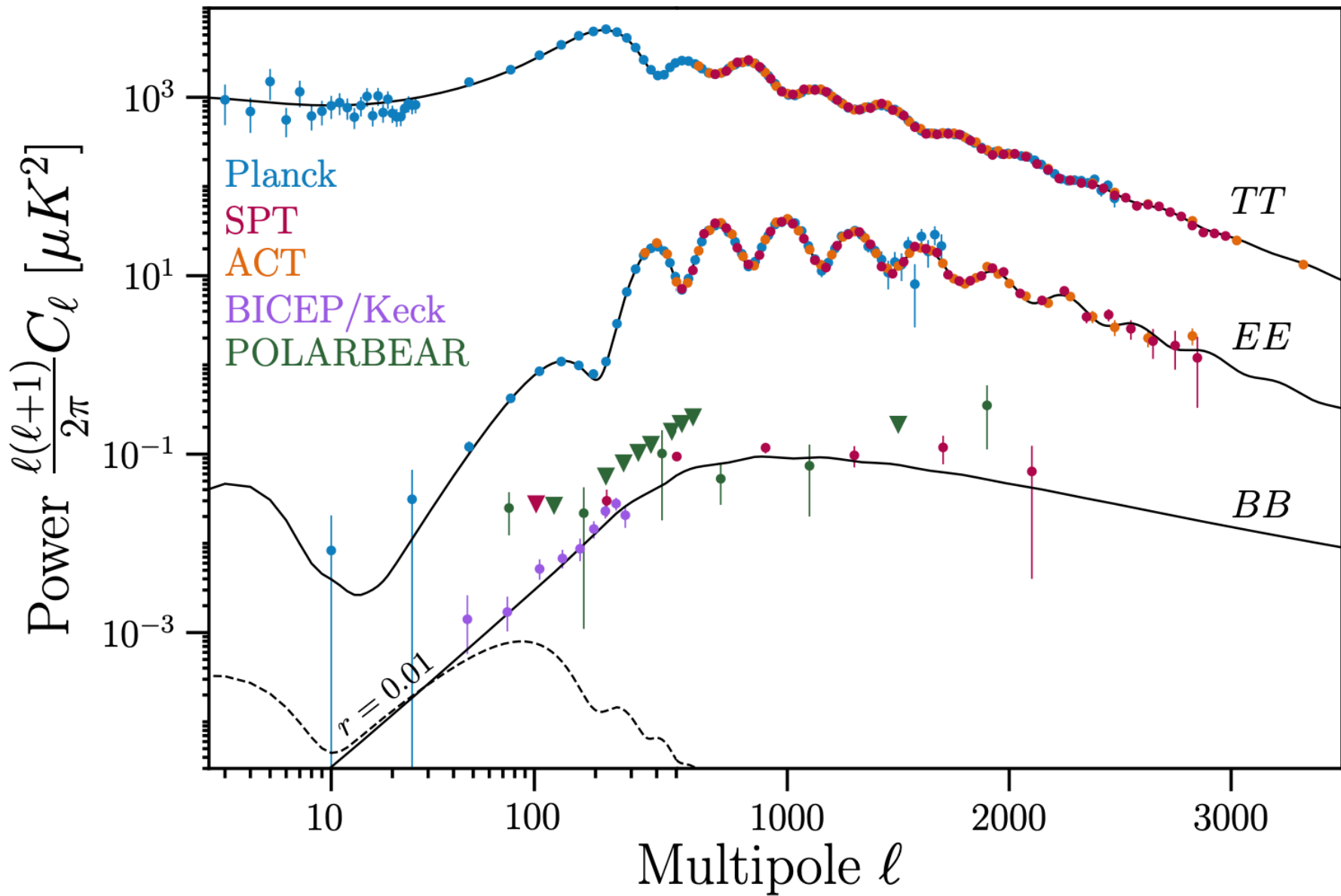
- CMB lensing power measures projected power of all matter (**no b**)
- galaxy clustering measures projected power of biased tracers (**b dN/dz**)²
- CMB lensing X galaxies measures projected power in common (**b dN/dz**)

Example: WISE X Planck lensing



Example: WISE X Planck lensing





Power spectrum Uncertainties

- fundamentally limited by number of independent measurements, noise
- $C_{l;\text{meas}} = C_{l;\text{true}} + C_{l;\text{noise}}$ *in any single map you can't tell the difference*
- $\text{Var}(C_l) \sim (2/n_{\text{meas}})C_l^2$ **“sample variance”**
- more modes means better measurement of $C_{l;\text{true}} + C_{l;\text{noise}}$
- lower noise gives better measure of $C_{l;\text{true}}$

Delensing lowers sample variance for B-mode searches

SPT-3G + external tracers
(galaxies+CIB) can remove
80% of lensing power

BICEP/Keck is signal-
dominated, so
delensing directly reduces the
error bar for constraints on
tensors
(also true for SPT-3G, for however
low in l can be reached)

